

Final Draft of the original manuscript

Nonnnenmacher, M.; Greenberg, D.:

Learning Implicit PDE Integration with Linear Implicit Layers.

In: The Symbiosis of Deep Learning and Differential Equations; NeuRIPS. 2021.

Learning Implicit PDE Integration with Linear Implicit Layers

Marcel Nonnenmacher

Institute of Coastal Systems, Helmholtz-Zentrum Hereon Geesthacht, Germany marcel.nonnenmacher@hereon.de

David S. Greenberg

Institute of Coastal Systems, Helmholtz-Zentrum Hereon Geesthacht, Germany david.greenberg@hereon.de

Abstract

Neural networks can learn local interactions to faithfully reproduce large-scale dynamics in important physical systems. Trained on PDE integrations or noisy observations, these emulators can assimilate data, tune parameters and learn subgrid process representations. However, implicit integration schemes cannot be expressed as local feedforward computations. We therefore introduce linear implicit layers (LILs), which learn and solve linear systems with locally computed coefficients. LILs use diagonal dominance to ensure parallel solver convergence and support efficient backward mode differentiation. As a challenging test case, we train emulators on semi-implicit integration of 2D shallow-water equations with closed boundaries. LIL networks learned compact representations of the local interactions controlling the 30.000 degrees of freedom of this discretized system of PDEs. This enabled accurate and stable LIL-based emulation over many time steps where feedforward networks failed.

1 Introduction

Numerical integration of partial differential equations (PDEs) is essential for many quantitative sciences, including fluid dynamics, numerical weather prediction and climate science [1, 6, 22]. Simulating these systems at scale requires substantial effort and computational resources [8, 36].

Recent studies used machine learning to emulate the PDE integration schemes that define full simulation models [12, 37, 29, 35] or their physical subcomponents [38, 26]. Applications include accelerated integration [37], system identification [18], automatic differentiation for data assimilation [2, 9, 24] and model tuning to match observations [32]. Many PDE integration schemes based on machine learning employ convolutional neural networks (CNNs) [18, 23, 21], whose kernels can precisely express spatial differential operators discretized as translation-invariant stencils.

However, a CNN that can express all of a PDE's mathematical terms may still fail to emulate numerical integration with implicit schemes, as required for stiff PDEs such as incompressible flow [1] and chemical dynamics [30]. Implicit schemes update the state by solving a systems of equations, and while the equation coefficients usually depend on local information, their solutions can depend on the entire previous state (Fig. 1a). This long-range spatial dependence prevents purely feed-forward networks from scaling efficiently to large systems.

To address this limitation, we introduce linear implicit layers (LILs), which define and solve a linear system with locally determined coefficients. This operation defines a nonlinear, differentiable implicit function with long-range spatial dependencies. LILs differ fundamentally from existing implicit layers based on Lipschitz constraints [7, 20], quadratic programming [39] or root-finding [5], instead constraining the linear system to be diagonally dominant to ensure convergence with a parallel solver in the forward pass. An efficient backward pass is implemented as a second linear solve. Using a challenging test problem based on the 2D shallow-water equations [34], we show that neural architectures incorporating LILs can learn dynamics integrated with semi-implicit schemes.

2 Methods

2.1 Problem Statement: Dynamical System Emulation

Let $x_0, x_{\Delta t}, x_{2\Delta t}, \ldots$ be a sequence of system states produced by numerical integration of a spaceand time-discretized PDE.¹ Given a collection of such sequences, our aim is to learn the update function $x_{t+\Delta t} = f(x_t)$.

From a data-centric perspective, f predicts the next sequence element. From a PDE-centric view, f numerically integrates over Δt time units, with the given discretization and integration schemes. As function family for f, we consider deep neural networks consisting of convolutional layers and LILs.

2.2 Linear implicit layers

Implicit layers can be written as $F_{\theta}(y, z) = 0$, where y is the layer input, θ the parameter vector and and z the layer output [19, 39]. We here consider the particular case of a system of linear equations,

$$A_{\theta}(y) \operatorname{vec}(z) = b_{\theta}(y). \tag{1}$$

where coefficients A_{θ} , b_{θ} are local functions of y and are parametrized by some vector θ . While the output z is a solution to a linear system, it depends nonlinearly on the inputs y. Firstly, since $A_{\theta}(y)$ and $b_{\theta}(y)$ are in general nonlinear functions, and second since $\text{vec}(z) = A_{\theta}(y)^{-1}b_{\theta}(y)$ depends bilinearly on $A_{\theta}(y)^{-1}$ and $b_{\theta}(y)$. We parameterize A_{θ} instead of its inverse to efficiently describe linear constraints z which are sparse and local, which would not be the case for $A_{\theta}(y)^{-1}$ (Fig. 1a).

When viewed as function $f_{\theta}(y) = z$, then for use as a network layer, we need to be able to evaluate f_{θ} and we need its partial derivatives wrt. A, b, θ and y for automatic differentiation. For evaluation of f_{θ} , all we need is a solver for systems of linear equations with matrix A_{θ} . For computing gradients of a target loss $\mathcal{L}_{\theta}(y)$ within reverse-mode differentiation, we need to solve another linear system with matrix A_{θ}^{\top} , as stated by straight-forward application of the implicit function theorem [17]:

$$A_{\theta}^{\top} \frac{\partial \mathcal{L}_{\theta}}{\partial \text{vec}(b_{\theta})} = \frac{\partial \mathcal{L}_{\theta}}{\partial \text{vec}(z)},\tag{2}$$

where $\frac{\partial \mathcal{L}}{\partial z}$ is the backpropagating gradient up to z. The implicit function theorem also gives $\frac{\partial \mathcal{L}}{\partial A} = -\frac{\partial \mathcal{L}}{\partial b} \text{vec}(z)^{\top}$, and the remaining partial derivatives for computing $\frac{\partial \mathcal{L}}{\partial y}$, $\frac{\partial \mathcal{L}}{\partial \theta}$ then are handled entirely by backpropagation. In particular, $\frac{\partial A}{\partial y}$, $\frac{\partial A}{\partial \theta}$ and $\frac{\partial b}{\partial y}$, $\frac{\partial b}{\partial \theta}$ depend on our choice for parametrizing the coefficients $A_{\theta}(y)$, $b_{\theta}(y)$.

2.3 Parametrization of the system of linear equations

For many applications $\dim(x_t)$ can easily reach between tens of thousands and billions [15]. This requires a sparse parametrization of the coefficient matrix A_{θ} , whose number of entries otherwise scales quadratically with system size. Importantly, A_{θ} also has to be full rank for the mapping $f_{\theta}(y) = z$ to be well-defined. We choose A_{θ} to be a banded matrix. To exploit translation equivariance common to PDEs, we compute the bands of $A_{\theta}(y)$ as convolutional outputs:

$$\operatorname{diag}_{\sigma_c}(A_{\theta}) = (\mathbf{M}_{\theta}(y))_c, \ c = 1, \dots, C, \tag{3}$$

where $M_{\theta}(y)$ is the output tensor of a CNN, c is a channel index, and $\operatorname{diag}_{\sigma}(A)$ indexes the main diagonal ($\sigma = 0$) and offdiagonals ($\sigma \neq 0$) of A. We choose the non-zero bands of A_{θ} to correspond

¹We focus on regular Cartesian grids, though our approach is not limited to these [10, 14].

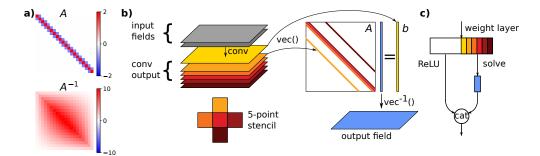


Figure 1: a) Sparse banded matrices representing convolutional stencils (upper: second derivative filter [-1,2,-1]) can have dense inverses (lower). b) The linear implicit layer generates output z (blue) by solving a linear system Az = b, with coefficients computed by convolutional layers applied to the inputs. c) Linear implicit layers are straightforward to embed in convolutional networks, by computing A, b from a subset of convolution outputs and concatenating the solution z with the rest.

to elements of computational stencils applied to output z: For systems with two spatial dimensions, we select the main diagonal and four off-diagonal bands (C=5) of A_{θ} to represent the central and outer points of a classical 5-point stencil, respectively (Fig. 1b). The righthand-side

$$b_{\theta} = (\mathbf{M}_{\theta}(y))_{c=C+1} \tag{4}$$

of the system of linear equations is computed as another convolutional output channel. By this choice of $A_{\theta}(y), b_{\theta}(y)$, the partial derivatives $\frac{\partial A}{\partial y}, \frac{\partial A}{\partial \theta}, \frac{\partial b}{\partial y}, \frac{\partial b}{\partial \theta}$ are computed from those of the network output $M_{\theta}(y)$ and, in the case of $\frac{\partial A}{\partial y}, \frac{\partial A}{\partial \theta}$, a simple look-up operation for where the bands are within A_{θ} . To avoid degeneracies, we normalize each row of A_{θ} and b_{θ} such that $\mathrm{diag}_0(A_{\theta})=1$.

2.4 Linear Solver

The main computational burden of linear implicit layers lies with the solver for the system of equations, which is called twice per gradient step. We use Red-Black Gauss-Seidel [28], an iterative solver that parallelizes well across large system sizes and only requires repeated (sparse) products of the form A_{θ} vec(z). This solver requires A_{θ} to be diagonally dominant, which we enforce through

$$\operatorname{diag}_{0}(A_{\theta}) = \sum_{\sigma \neq 0} |\operatorname{diag}_{\sigma}(A_{\theta})| + \exp((\mathbf{M}_{\theta})_{k}), \tag{5}$$

where $\sigma_k = 0$, before normalizing rows of A_{θ} , b_{θ} . Diagonal dominance ensures A_{θ} is full rank.

2.5 Boundary conditions

A particular challenge [23] for the definition of the system of equations is posed by boundary conditions (BCs), which enforce equations on the domain boundaries Ω that can differ drastically from those in the interior. With binary boundary map I_{Ω} we compose the full system of equations as

$$\mathbf{M}_{com}(y) = (1 - I_{\Omega}) \,\mathbf{M}_{\theta}(y) + I_{\Omega} \,\mathbf{M}_{BC}(y), \tag{6}$$

and analogous for b_{com} . The equations on the boundary are defined by I_{Ω} M_{BC} , I_{Ω} b_{BC} , which we can learn with a dedicated model, or in simple cases define manually.

3 Numerical experiments on Shallow-Water Equations

We test² our implicit integration models on the shallow-water equations (SWE) in two spatial dimensions, with bottom friction and variable depth profile h. The SWE describe the interactions of water height ξ , vertical and horizontal velocities $v = (v_x, v_y)$. We use a semi-implicit integration

²Code for reproducing our results is found at https://github.com/m-dml/lil2021swe

scheme [4] that advances the system state $x_t = (\xi_t, v_t, h)$ by Δt in three steps:

$$v^* = f^*(\xi_t, v_t, h)$$
 interim velocities (explicit) (7)

$$A^*(v^*) \xi_{t+\Delta t} = b^*(v^*) \qquad \text{implicit treatment of water height}$$

$$v_{t+\Delta t} = f^{\Delta}(\xi_{t+\Delta t}, v^*) \qquad \text{explicit velocity update}$$

$$(9)$$

$$v_{t+\Delta t} = f^{\Delta}(\xi_{t+\Delta t}, v^*)$$
 explicit velocity update (9)

We use an Arakawa C-grid [3] with 100x100 grid points at 10km resolution (Fig. 2a). Depth profile maps h are drawn from correlated Gaussian noise, and BCs enforce $\xi, v_x, v_y = 0$ on the boundary.

Models are trained on 60.000 numerically integrated pairs $(x_t, x_{t+\Delta t})$, with a mean-squared loss to predict the next system state, and data variances normalized per channel. The data consists of numerical simulations from different initial conditions and depth profiles. Initial conditions are $\xi, v_x, v_y = 0$ with different local perturbations to water height ξ of varying strength, spatial extent and location. For each initial condition we simulate 25 hours of dynamics. We use 80% of simulations for training and 20% as validation data. For model training we use Adam [16] with default parameters. We train 5-layer CNNs with 3x3 kernels, both ReLU and bilinear activations [11] and 14.4k free parameters. Six output channels of the 4th convolutional layer are used to define $M_{\theta}(x_t)$ and $b_{\theta}(x_t)$ for a 5-point stencil system A_{θ} $\hat{\xi}_{t+\Delta t}=b_{\theta}$. $\hat{\xi}_{t+\Delta t}$ is passed to the final network layer to estimate $\hat{v}_{t+\Delta t}$. For comparison, we train fully explicit models (pure CNNs) in which the channels for M_{θ} and b_{θ} act as regular convolutional channels. We enforce BCs on ξ for the semi-implicit models at the level of M_{com} , b_{com} , and for the explicit models and v we enforce them directly on the model output through masking with I_{Ω} .

We train the models on three datasets with different integration step-sizes $\Delta t \in \{300s, 900s, 1500s\}$. All experiments combined required 160 hours on a V100 GPU. For evaluation we compare trained models to numerical integration on novel initial conditions. Models with LILs perform well for all step sizes, while pure CNNs fail at high Δt (Fig. 2b-c), for which long-range dependencies emerge.

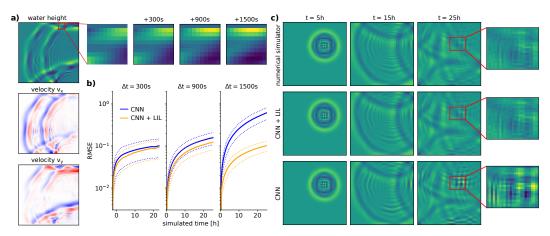


Figure 2: a) Shallow water equation system state (left) and magnified surface height dynamics (right). b) Mismatch between numerical integration and emulation with CNNs with and without LILs vs. simulated time, for different Δt . RMSE ± 1 std over 100 initial conditions. c) Comparison of CNNs with and without LILs over 20h with $\Delta t = 1500$ s, starting from the same state 5h after initial perturbation.

Discussion

We showed that LILs enable deep learning of spatially extended dynamics that require semi-implicit integration. Previously Look et al. [19] studied implicitly-integrated neural models on ordinary differential equations with up to 150-dimensional system states. Here we focused on scalability in PDE systems with multiple spatial axes and variable fields, reaching 10.000 grid points with 3 variables each. Choosing diagonally-dominant banded matrices allows fast parallel solvers to scale to large systems, but also limits the expressiveness of LILs. For instance, LILs could learn a fully implicit model scheme for the SWE by defining a system of equations over all system variables, but with larger systems the requirement of diagonal dominance becomes increasingly restricting, especially for large time steps.

Other LIL parametrizations such as low-rank perturbations of diagonal matrices are possible, as would be other solvers such as conjugate gradient methods [31]. While we used LILs to incorporate the long-range input output dependence appearing in (semi-)implicit schemes, other approaches are possible as well, such as U-Nets [27, 35] and self-attention [33]. Unlike LILs, U-Nets have a restricted dependency range by construction, and self-attention for multi-axis systems typically is spatially restricted as well [25, 13].

LILs provide efficient and straightforward backward-mode differentiation, and require relatively few parameters, requiring only six convolutional channels per output field with a 5-point stencil. Several LILs can be stacked within a network. These features suggest that LILs could prove useful as general building blocks for deep models requiring non-local computations.

5 Acknowledgments

M. Nonnenmacher and D. Greenberg were supported by the Helmholtz AI initiative. We thank Kai Logemann for help with the SWE numerical simulator, and Jan-Matthis Lueckmann, Tobias Machnitzki, Shivani Sharma and Vadim Zinchenko for comments on the manuscript.

References

- [1] John David Anderson and J Wendt. Computational fluid dynamics, volume 206. Springer, 1995.
- [2] A Apte, Christopher KRT Jones, AM Stuart, and Jochen Voss. Data assimilation: Mathematical and statistical perspectives. *International journal for numerical methods in fluids*, 56(8):1033–1046, 2008.
- [3] Akio Arakawa and Vivian R Lamb. Computational design of the basic dynamical processes of the ucla general circulation model. *General circulation models of the atmosphere*, 17(Supplement C):173–265, 1977.
- [4] Jan O Backhaus. A semi-implicit scheme for the shallow water equations for application to shelf sea modelling. *Continental Shelf Research*, 2(4):243–254, 1983.
- [5] Shaojie Bai, J Zico Kolter, and Vladlen Koltun. Deep equilibrium models. *arXiv preprint* arXiv:1909.01377, 2019.
- [6] Peter Bauer, Alan Thorpe, and Gilbert Brunet. The quiet revolution of numerical weather prediction. *Nature*, 525(7567):47–55, 2015.
- [7] Jens Behrmann, Will Grathwohl, Ricky T. Q. Chen, David Duvenaud, and Jörn-Henrik Jacobsen. Invertible Residual Networks. *arXiv:1811.00995* [cs, stat], May 2019. arXiv: 1811.00995.
- [8] Stewart Cant. High-performance computing in computational fluid dynamics: progress and challenges. *Philosophical Transactions of the Royal Society of London. Series A: Mathematical, Physical and Engineering Sciences*, 360(1795):1211–1225, 2002.
- [9] Alberto Carrassi, Marc Bocquet, Laurent Bertino, and Geir Evensen. Data assimilation in the geosciences: An overview of methods, issues, and perspectives. *Wiley Interdisciplinary Reviews: Climate Change*, 9(5):e535, 2018.
- [10] Taco S Cohen, Mario Geiger, Jonas Köhler, and Max Welling. Spherical cnns. *arXiv preprint* arXiv:1801.10130, 2018.
- [11] Ronan Fablet, Said Ouala, and Cédric Herzet. Bilinear residual neural network for the identification and forecasting of geophysical dynamics. In 2018 26th European Signal Processing Conference (EUSIPCO), pages 1477–1481. IEEE, 2018.
- [12] Radek Grzeszczuk, Demetri Terzopoulos, and Geoffrey Hinton. Neuroanimator: Fast neural network emulation and control of physics-based models. In *Proceedings of the 25th annual conference on Computer graphics and interactive techniques*, pages 9–20, 1998.

- [13] Jonathan Ho, Nal Kalchbrenner, Dirk Weissenborn, and Tim Salimans. Axial attention in multidimensional transformers. *arXiv preprint arXiv:1912.12180*, 2019.
- [14] Chiyu Jiang, Jingwei Huang, Karthik Kashinath, Philip Marcus, Matthias Niessner, et al. Spherical cnns on unstructured grids. *arXiv preprint arXiv:1901.02039*, 2019.
- [15] Hyun-Gyu Kang, Katherine J Evans, Mark R Petersen, Philip W Jones, and Siddhartha Bishnu. A scalable semi-implicit barotropic mode solver for the mpas-ocean. *Journal of Advances in Modeling Earth Systems*, 13(4):e2020MS002238, 2021.
- [16] Diederik P Kingma and Jimmy Ba. Adam: A method for stochastic optimization. arXiv preprint arXiv:1412.6980, 2014.
- [17] Steven G Krantz and Harold R Parks. *The implicit function theorem: history, theory, and applications.* Springer Science & Business Media, 2012.
- [18] Zichao Long, Yiping Lu, Xianzhong Ma, and Bin Dong. Pde-net: Learning pdes from data. In International Conference on Machine Learning, pages 3208–3216. PMLR, 2018.
- [19] Andreas Look, Simona Doneva, Melih Kandemir, Rainer Gemulla, and Jan Peters. Differentiable implicit layers. *arXiv preprint arXiv:2010.07078*, 2020.
- [20] Cheng Lu, Jianfei Chen, Chongxuan Li, Qiuhao Wang, and Jun Zhu. Implicit normalizing flows. *arXiv preprint arXiv:2103.09527*, 2021.
- [21] Suryanarayana Maddu, Dominik Sturm, Bevan L Cheeseman, Christian L Müller, and Ivo F Sbalzarini. Stencil-net: Data-driven solution-adaptive discretization of partial differential equations. *arXiv* preprint arXiv:2101.06182, 2021.
- [22] K McGuffie and A Henderson-Sellers. Forty years of numerical climate modelling. *International Journal of Climatology: A Journal of the Royal Meteorological Society*, 21(9):1067–1109, 2001.
- [23] Arvind T Mohan, Nicholas Lubbers, Daniel Livescu, and Michael Chertkov. Embedding hard physical constraints in convolutional neural networks for 3d turbulence. In *ICLR 2020 Workshop on Integration of Deep Neural Models and Differential Equations*, 2020.
- [24] Marcel Nonnenmacher and David S Greenberg. Deep emulators for differentiation, forecasting and parametrization in earth science simulators. *Journal of Advances in Modeling Earth Systems*, 13(7), 2021.
- [25] Niki Parmar, Ashish Vaswani, Jakob Uszkoreit, Lukasz Kaiser, Noam Shazeer, Alexander Ku, and Dustin Tran. Image transformer. In *International Conference on Machine Learning*, pages 4055–4064. PMLR, 2018.
- [26] Ali Ramadhan, John Marshall, Andre Souza, Gregory LeClaire Wagner, Manvitha Ponnapati, and Christopher Rackauckas. Capturing missing physics in climate model parameterizations using neural differential equations. *arXiv preprint arXiv:2010.12559*, 2020.
- [27] Olaf Ronneberger, Philipp Fischer, and Thomas Brox. U-net: Convolutional networks for biomedical image segmentation. In *International Conference on Medical image computing and computer-assisted intervention*, pages 234–241. Springer, 2015.
- [28] Yousef Saad. Iterative methods for sparse linear systems. SIAM, 2003.
- [29] Alvaro Sanchez-Gonzalez, Jonathan Godwin, Tobias Pfaff, Rex Ying, Jure Leskovec, and Peter W. Battaglia. Learning to Simulate Complex Physics with Graph Networks. *arXiv* preprint arXiv:2002.09405 [physics, stat], 2020.
- [30] Adrian Sandu and Rolf Sander. Simulating chemical systems in fortran90 and matlab with the kinetic preprocessor kpp-2.1. *Atmospheric Chemistry and Physics*, 6(1):187–195, 2006.
- [31] Gerard LG Sleijpen and Diederik R Fokkema. Bicgstab (ell) for linear equations involving unsymmetric matrices with complex spectrum. *Electronic Transactions on Numerical Analysis.*, 1:11–32, 1993.

- [32] Kiwon Um, Philipp Holl, Robert Brand, and Nils Thuerey. Solver-in-the-loop: Learning from differentiable physics to interact with iterative pde-solvers. *arXiv preprint arXiv:2007.00016*, 2020.
- [33] Ashish Vaswani, Noam Shazeer, Niki Parmar, Jakob Uszkoreit, Llion Jones, Aidan N Gomez, Łukasz Kaiser, and Illia Polosukhin. Attention is all you need. In *Advances in neural information processing systems*, pages 5998–6008, 2017.
- [34] Cornelis Boudewijn Vreugdenhil. *Numerical methods for shallow-water flow*, volume 13. Springer Science & Business Media, 1994.
- [35] Nils Wandel, Michael Weinmann, and Reinhard Klein. Unsupervised deep learning of incompressible fluid dynamics. *arXiv preprint arXiv:2006.08762*, 2020.
- [36] Warren M Washington, Lawrence Buja, and Anthony Craig. The computational future for climate and earth system models: on the path to petaflop and beyond. *Philosophical Transactions of the Royal Society A: Mathematical, Physical and Engineering Sciences*, 367(1890):833–846, 2009.
- [37] Steffen Wiewel, Moritz Becher, and Nils Thuerey. Latent space physics: Towards learning the temporal evolution of fluid flow. *Computer graphics forum*, 38(2):71–82, 2019.
- [38] Janni Yuval and Paul A O'Gorman. Stable machine-learning parameterization of subgrid processes for climate modeling at a range of resolutions. *Nature communications*, 11(1):1–10, 2020.
- [39] Qianggong Zhang, Yanyang Gu, Michalkiewicz Mateusz, Mahsa Baktashmotlagh, and Anders Eriksson. Implicitly Defined Layers in Neural Networks. *arXiv:2003.01822 [cs]*, March 2020. arXiv: 2003.01822 version: 1.