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Sediment Dynamics in the Black Sea: Numerical Modelling and Remote Sensing Observations

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Abstract

Here, we address the sediment dynamics in the Black Sea based on analysis of remote sensing data from MERIS satellite spectrometer and numerical simulations with NEMO model. Boundary conditions consist of realistic meteorological forcing, including significant wave height generated by WAM model. A number of sensitivity runs was analysed with the aim to find the most suitable parameters governing sediment fluxes. The comparison between numerical simulations and remote sensing data gives credibility to the quality of simulations. The combined effect of wind waves and currents in the bed layer controls the sediment resuspension that appears to be the major basin wide source of sediment. Sensitivity experiments included or excluded different forcing terms, e. g. sediment flux from rivers enable to determine the spatial extensions of different point-sources. It is concluded that wind-wave forcing is manifested in the sediment dynamics through episodic high energy events contributing to the increase of horizontal sediment fluxes over the north-western shelf. Both satellite images and numerical model simulations demonstrated that the penetration of suspended sediment into the basin interior is governed by the dynamics of coastal and open-ocean eddies. While fine sediment at sea surface could cross the continental slope extending into the open ocean, coarser fractions follow the bottom and their penetration into the open ocean is limited. The conclusion is thus that the deposition patterns correlate with the specific shape of Black Sea topography largest depositions are observed in the area of continental slope. Analysis of numerical simulations could motivate further observations aiming to verify the possibility identified in the numerical simulations that winter convection/mixing could also affect the propagation of fine suspended sediment.

1 Introduction

The Black Sea is a nearly enclosed deep basin connected to the Sea of Marmara and the Sea of Azov by the narrow Bosphorus and Kerch Straits, respectively (figure 1). The area shallower than 200 m (light blue area, which gives 27 percent of the total area), is mainly located in the northwest. Black Sea catchment area of about 1,864,000 km² covers large parts of Europe and Asia (figure 1), providing a total freshwater supply of 3×10^2 km³ per year. This large amount of fresh water flux along with the narrow opening in the strait of Bosphorus make the Black Sea a typical estuarine basin where the surface salinity is about half that of the Mediterranean's. The strong vertical (haline) stratification limits the vertical exchange and creates a unique chemical and biological environment.

The circulation in the Black Sea is structured in two connected gyre systems called Eastern and Western Gyres encompassing the basin (the Rim Current, figure 2). This is a frontal current associated with a difference of ~ 0.2 m between sea level in the coastal and open sea. Its position is controlled by the bottom topography, which is characterised by a gentle slope in the northwestern and western Black Sea and abrupt one along the southern and eastern coasts (see the contrast between light blue and dark blue areas in figure 1). Anticyclonic eddies are formed between the Rim Current and the coast, the most prominent of them are the Batumi and Sevastopol eddies (figure 2)

Annual production of organic mater of 818×10^6 t yr⁻¹ exceeds substantially the total sediment load from the rivers around the Black Sea periphery of about 150×10^6 t yr⁻¹ (figure 3, see also Shimkus and Trimonis, 1974; Panin and Jipa, 2002; Balkas et al., 1990; Bondar et al., 1991; Panin, 1996). About 65 percent of this amount enters into the northwestern shelf region (Hay, 1994). The sedimentary and ecological systems of the continental shelf are dominated by the fluxes from the major rivers, the Danube, the Dniester, and the Dnieper (Oguz, 2005). The River Danube, which accounts for ~ 60 percent of the total river input into the northwestern continental shelf discharges 45 to 50×10^6 t yr⁻¹ of suspended solids into the Black Sea (Popa 1993). The Danube Delta is subjected to waves primarily approaching from the east, driving net sediment transport towards the southwest (Giosan et al., 1999; Giosan et al., 2005). Under some specific circulation conditions (Stanev et al., 2002) the

57 Danube plume is entrapped by an anticyclonic shelf-gyre extending to the north.
58 Average sedimentation rates for the Upper Quaternary range between 1.19 and 2.19
59 m ka^{-1} for the Danube fan and between 1.07 and 2.03 m ka^{-1} for the Dnieper fan
60 (Winguth et al., 2000). These values are much higher than in other Mediterranean
61 basins, which is caused by the large drainage area of the Black Sea. The thickness
62 of sediment layers that have accumulated over the last 3000 years of the Black Sea
63 history is 20 to 80 cm depending on the region. Sediment trap data (Hay and Honjo,
64 1989) reveal large variations in the total particle flux (from 2 to 71 $\text{mg m}^{-2}\text{day}^{-1}$).
65 Fine-grained terrigenous matter may remain in suspension in the surface water for
66 several months, which may be a sufficient time to transport the suspended matter
67 across the basin before it settles to the sea floor.

68 The sediment pattern in the near-shore zone of the Black Sea is governed by the
69 surface and alongshore bottom currents and wave action (Ross et al., 1978; Nikolov
70 et al., 2006). In the western Black Sea suspended particulate mater (SPM) from the
71 Danube, Dnieper, and other big rivers is deposited and trapped on the broad western
72 shelf. In the eastern basin, the terrigenous material easily crosses the narrow shelf
73 and enters the deep basin.

74 SPM is also referred to as Total Suspended Mater (TSM) in the literature and we will
75 use here the two abbreviations as synonyms. SPM in the Black Sea has terrestrial,
76 planktonic, bacterial and fossil origin. Grain size fraction $< 2\mu\text{m}$ is characteristic for
77 the abyssal plane, 6.3-20 μm for the slope area and 20 – 63 μm for the shelf, but also
78 for the slope area (Müller and Stoffers, 1974). According to Wakeham et al. (1991)
79 the vertical profile of TSM displays high concentrations in surface waters ($\sim 400\mu\text{g/l}$
80 at 2 m), which decreases to $\sim 150\mu\text{g/l}$ at 200 m and does not change much down
81 to 1000 m (figure 4). Noteworthy is that TSM gives the overall picture of sediment
82 distribution, which includes also plankton ($< 63\mu\text{m}$), the later is dominant in the
83 interior Black Sea. The separation between TSM of different origin is a difficult
84 issue, also when one analyses satellite data (see further in the text).

85 Recent remote sensing data contributed to further understanding of the basin-wide
86 sediment dynamics and provided valuable material for the validation of numerical
87 models. Sur et al. (1996) used satellite (CZCS, AVHRR) and in-situ (CTD, ADCP)
88 data to characterise the impact of meso-scale motions and boundary currents on
89 the transport and productivity along the wide Western Black Sea shelf. They found

90 various forms of isolated transient features such as filaments, coherent dipole and
91 monopole eddies, and jets to contribute substantially to the transport of SPM across
92 the continental shelf.

93 Karabashev et al. (2006) demonstrated the potential of SeaWiFS imagery (figure 5)
94 to describe the horizontal patterns (vortex dipoles) associated with the coastal/deep
95 basin water exchange. Their remote sensing data indicated that the patterns are
96 influenced by sediments from the Danube Delta. Similar patches and filaments were
97 observed earlier in the same area by Ginzburg et al. (1997) and Guneroglu et al.
98 (2005).

99 Kopelevich et al. (2004) analysed the spatial and temporal variability of the bio-
100 optical characteristics in the Black Sea derived from SeaWiFS data in 1998 - 2001
101 and concluded that fluviially derived particles are spread over the entire basin. The
102 particulate matter brought by the rivers in the northwestern part (mainly by the
103 Danube River) is transported along the coast by the Rim Current and across the
104 shelf-continental slope area by meso-scale eddies and via turbulent exchange (see,
105 as an example, their Fig. 12).

106 Unlike to observations, numerical modelling of sediment transport in the Black Sea
107 is still underdeveloped, differently from the case of the neighboring Mediterranean
108 Sea where several good examples already exist (Wang and Pinardi, 2002; Kourafalou
109 et al., 2004; Sherwood et al., 2004; Ulses et al., 2008; Harris et al., 2008, Bevera
110 et al., 2009) demonstrating very promising results. It has been established in these
111 studies that the wave-induced remobilisation of the deposits is a key mechanism of
112 cross-shelf sediment transport (see also Traykovski et al., 2000). To our knowledge,
113 the only experience existing in the field of numerical modelling of (coarse) sediment
114 in the Black Sea is due to Giosan et al. (2005), who address the morphodynamics
115 of the wave-influenced Danube Delta.

116 The motivation of this study is to initiate first steps toward development of basin
117 scale modelling capabilities for sediment transport in the Black Sea. Because we
118 address the entire basin, we include in our numerical simulations not only coarse,
119 but also fine sediment fractions which could propagate far from their river and
120 bottom origin and sometimes penetrate the basin interior. Evaluating geomorphic
121 evolution or delta growth is not addressed in the present study, neither the bed-load
122 transport. Dealing with fine sediment leads to a substantial difficulty, which is the

123 separation between sediment from rivers and sediment remobilized from bottom,
124 from one side, from the sediment due to biological productivity, from the other.
125 Modelling approach to this problem would necessitate using coupled biological and
126 sediment models with an appropriate description of the processes associated with
127 the interaction of these two kinds of sediment and grain size variation. Further
128 in this direction is the issue of simulating optical properties affecting biological
129 productivity. The present study, which could be considered as a first step, does not
130 consider the biogeochemistry, therefore the research questions addressed are limited
131 to physical processes and aim to: (1) present a sediment transport model for the
132 Black Sea, (2) demonstrate the realism of simulations, mainly through comparison
133 with remote sensing data, (3) learn more about the Black Sea sediment dynamics.
134 One additional motivation of this study is to identify some important problems in
135 sediment modelling in the Black Sea and, based on solved and unsolved problems,
136 to formulate the direction of future steps in this field. Following these lines we first
137 describe the numerical model in section 2, followed by the analysis of results in
138 section 3 and conclusions.

139 2 Models and data

140 2.1 Brief description of the model

141 Circulation model

142 Extended references on numerical modelling of the Black Sea circulation are provided
143 in several reviews (Stanev, 1990; Stanev et al., 2002; Kara et al., 2005, Stanev, 2005).
144 Although there are many models used so far, one relatively new one Nucleus for
145 European Modelling of the Ocean (NEMO) (Madec, 2008) deserves our attention,
146 therefore, we describe here very briefly its setup for the Black Sea (for more details
147 see Grayek et al., 2010). The appropriate resolution of important physical processes,
148 rich physical parametrisations, and conservation properties justify our selection of
149 this model.
150 The NEMO is discretized on an Arakawa C grid. The horizontal resolution in
151 our Black Sea setup is $1/12^\circ$, thus, we have a matrix of 133×76 points, which is
152 the same as in our similar previous works based on the Modular Ocean Model

(MOM) setup for the Black Sea (Stanev et al., 2003). The investigated area extends from $27^{\circ}38'E \times 40^{\circ}45'N$ to $42^{\circ}05'E \times 46^{\circ}70'N$. The vertical grid uses a hard-wired hyperbolic tangent stretching function with 31 levels. Our setup employs also partial step representation of bottom topography (Barnier et al., 2006) in order to adjust the model depth to the observed one.

For lateral mixing of tracers (temperature, salinity and sediment) we use a second-order geopotential operator (Laplacian) with $A^{lT} = 100\text{m}^2\text{s}^{-1}$, where indices "l" and "T" stay for "lateral" and "tracer", correspondingly. A geopotential bi-laplacian operator (Griffies, 2004) parametrizes the horizontal eddy viscosity with $A^{lm} = -7 \times 10^9\text{m}^4\text{s}^{-1}$, where "m" stays for "momentum". The vertical (index "v") eddy viscosity A^{vm} and diffusivity A^{vT} coefficients are computed from a closure scheme (Madec, 2008) with a prognostic equation for the turbulent kinetic energy \bar{e} and a closure assumption for the turbulent length scales, so that the time evolution of \bar{e} is a result of its production through vertical shear, destruction through stratification, vertical diffusion, and dissipation:

$$\begin{aligned} \frac{\partial \bar{e}}{\partial t} &= A^{vm} \left[\left(\frac{\partial u}{\partial z} \right)^2 + \left(\frac{\partial v}{\partial z} \right)^2 \right] - A^{vT} N^2 + \frac{\partial}{\partial z} \left[A^{vm} \frac{\partial \bar{e}}{\partial z} \right] - c_{\epsilon} \frac{\bar{e}^{3/2}}{l_{\epsilon}} \\ A^{vm} &= C_k l_k \sqrt{\bar{e}} \\ A^{vT} &= A^{vm} / P_{rt} \end{aligned} \quad (1)$$

where u and v are the horizontal velocity components, N is the local Brunt-Vaisala frequency, l_k and l_{ϵ} are the dissipation and mixing length scales, P_{rt} is the Prandtl number, $C_k = \sqrt{2}/2$ and $C_{\epsilon} = 0.1$. In addition, an enhanced vertical diffusion is applied on both tracer and momentum which means assigning large value $A_{max}^{vT} = A_{max}^{vm} = 1\text{m}^2\text{s}^{-1}$ to the vertical eddy diffusivity coefficients in regions with unstable stratification, whereas their minimal values are set to be $A_{min}^{vT} = 1.2 \times 10^{-5}\text{m}^2\text{s}^{-1}$ and $A_{min}^{vm} = 1.2 \times 10^{-4}\text{m}^2\text{s}^{-1}$. The selected advection scheme for tracers is the total variation diminishing (TVD) scheme, which is, at best, second-order accurate (Barnier et al., 2006). TVD scheme may be viewed as a nonlinear flux limiter that adds just enough diffusion to prevent numerical instabilities.

178 Sediment transport model

179 The sediment transport model uses a standard diffusion-advection equation (the
 180 same as for the temperature and salinity) for the SPM concentration c with an
 181 extra sinking term added on the right hand side:

$$\frac{\partial c}{\partial t} + u \frac{\partial c}{\partial x} + v \frac{\partial c}{\partial y} + w \frac{\partial c}{\partial z} = A^{lT} \Delta c + \frac{\partial}{\partial z} \left(A^{vT} \frac{\partial c}{\partial z} \right) + \frac{\partial}{\partial z} (w_s c) \quad (2)$$

182 where w is the vertical velocity component and w_s is the settling velocity of the
 183 sediments in suspension. Here we refer to these particles as to mud or fine SPM ($d <$
 184 $63\mu m$) and we assume that they are non-cohesive (do not flocculate or aggregate)
 185 and that their concentration (relatively small) has no effect on the water density.

186 Deposition and erosion

187 Deposition and erosion rates are controlled by the bed shear stress $\tau_b = A^{vm}(\partial u/\partial z)$.
 The sediment flux at the sea bed is

$$\left(A^{vT} \frac{\partial c}{\partial z} \right)_{z=H} = \begin{cases} -D & \tau_b < \tau_d \\ E & \tau_b > \tau_e \end{cases} \quad (3)$$

188 where D and E are deposition and erosion rates. Deposition rate given by Einstein
 189 and Krone (1962) is

$$D = w_s \times c_b \left(1 - \frac{\tau_b}{\tau_d} \right), \quad (4)$$

190 where c_b is the bottom concentration and τ_d is the critical shear stress for deposition.
 191 The definition of the bottom concentration c_b is as for the remaining (standard)
 192 tracers in NEMO, and uses the modifications associated with partial steps, connected
 193 with the fact that bottom depths do not equal the standard model depths (see
 194 Madec, 2008).

195 The erosion rate is computed using the formula of Partheniades (1965):

$$E = M_e \left(\frac{\tau_b}{\tau_e} - 1 \right), \quad (5)$$

196 where M_e is an empirical parameter giving the erosion rate at twice the critical shear
 197 stress for erosion τ_e .

198 The critical shear stress for deposition is chosen to be equal to that for erosion:
 199 $\tau_d = \tau_e = \tau_{crit}$. This means that either deposition or erosion occurs and no region of
 200 transition exists where neither of the two processes are active. Following previous
 201 studies Wang and Pinardi (2002) we assume $\tau_{crit} = 2 \times 10^{-2} Nm^{-2}$.

202 **2.2 Model forcing, boundary conditions and experiments**

203 The data used here to construct the forcing consists of radiation (including clouds),
 204 temperature and relative humidity at 2 m and atmospheric winds at 10m. They are
 205 produced by the European Centre for Medium-Range Weather Forecasts (ECMWF)
 206 and were kindly made available in the frame of SESAME (Southern European Seas:
 207 Assessing and Modelling Ecosystem changes) project by the INGV, Italy. The tem-
 208 poral sampling rate is 6 hours and the spatial resolution is $0.25^\circ \times 0.25^\circ$.

209 It is important to mention that the actual forcing in the model is derived interactively
 210 - that is the heat fluxes depend on the sea surface temperature (SST), which is a
 211 model variable. This type of forcing has been proved successful in regional ocean
 212 models of the Mediterranean (Roussenov et al, 1995) and the Black Sea (Staneva et
 213 al., 1995; Stanev et al., 2003).

214 A spin-up of one year forced with climatological data (mean year) has first been
 215 performed, followed by an integration for the period 2006-2009 with realistic forcing
 216 data. All experiments are forced with time-varying meteorological condition.

217 **Momentum flux**

218 The momentum flux in the model is introduced through the wind stress components:

$$(219) \quad (\tau_x, \tau_y) = \rho_a C_d |W| (W_x, W_y) \quad (6)$$

219 where (W_x, W_y) are the wind components at 10 m, ρ_a is the density of air and C_d
 220 is the drag coefficient, which is stability dependent (Grayek et al., 2010). This me-
 221 chanical forcing is shown in figure 6a as the annual mean wind. The area mean

222 wind-stress curl is positive throughout the year with largest positive anomalies (cy-
 223 clonic curl) located in the eastern Black Sea. Overall, the spatial pattern is consistent
 224 with the similar ones displayed by Stanev et al.(2003) and Grayek et al. (2010).

225 **Thermohaline fluxes**

226 Surface buoyancy forcing is controlled by thermohaline fluxes shaped by air temper-
 227 ature and humidity as well as precipitation. Evaporation is diagnosed in the model
 228 using bulk aerodynamic formulas. The total heat flux is

$$Q_T = Q_s - H - LE - Q_b, \quad (7)$$

229 where Q_s is the downward flux of solar radiation, H is the sensible heat flux, E is
 230 evaporation, L the latent heat of evaporation and Q_b is the upward flux resulting
 231 from the long wave radiation lost from the sea surface. In our model, we used the
 232 same parametrisations for the individual components of heat fluxes as in Stanev et
 233 al. (2003).

234 Virtual surface salinity flux is

$$Q_{sal} = Q_W S, \quad (8)$$

235 where $Q_W = R + P - E$ is the difference between the river runoff (R) and precipitation
 236 (P), from one side and the evaporation (E) from the other, S is the surface salinity
 237 (18 psu in the Black Sea). The approximation $S = 18psu$ in equation (8) is good
 238 except in the vicinity of Danube Delta. However, its advantage is that a globally
 239 balanced $P - E$ field converts to a globally balanced surface salinity flux (i.e., the
 240 global mean salinity in a model does not change). The salinity balance is closed in
 241 the strait of Bosphorus (see further in the text).

242 The model has been initialised with horizontally uniform temperature and salinity
 243 profiles taken from climatological data. The geothermal heat and salinity flux at
 244 the bottom were neglected.

245 **Fluxes from rivers and Bosphorus Strait**

246 The river runoff data used in the model (figure 7a) is based on the monthly mean
 247 data for the period 1960-2000 (Ludwig et al.,2009). The major coastal sources of

SPM are defined as point sources (river mouths) and the values are computed from the data presented by Panin and Jipa (2002, see their Table 1). The major non-point source (sink) of sediments is the sea floor.

The horizontal boundaries were closed for the Kerch Strait and open for the Bosphorus Straits. This formulation follows the consensus that the transport in the Bosphorus Straits is crucial for the Black Sea oceanography. As shown by Peneva et al. (2001), there is a clear correlation between the river runoff and the barotropic straits transport. Therefore, we estimated the barotropic transport in the Straits of Bosphorus R_{Bo} as a linear function of the river runoff R_{ri} . The monthly variation of the water outflow R_{Bo} was constructed according to the monthly mean data $R_{ri,L}$ of Ludwig et al. (2009) with "corrected" mean value relative to the one of the annual mean climatological river runoff $R_{ri,K}$ of Kara et al. (2007, see their table 7):

$$R_{Bo} = k \{ R_{ri,L} - (\overline{R_{ri,L}} - \overline{R_{ri,K}}) \} \quad (9)$$

where $\overline{(\cdot)}$ is the annual mean and the coefficient k is found using the least mean squares method for the function $R_{Bo}(R_{ri,K})$. The inflow from the Bosphorus Straits is applied in the model as "negative" river at a grid point near the shore (so that $R = R_{ri} - |R_{Bo}|$) whereas the corresponding salinity flux Q_{sal}^{Bo} is computed as in equation (8) with $S = 36psu$, which reflects Marmara Sea conditions.

Tidal processes were neglected because for the Black Sea (small and enclosed basin) tides have very small amplitude.

Bed shear stress

We represent the bed shear stress

$$\tau_b = \tau_b^c + \tau_b^w$$

as a sum of two terms: due to currents (index "c") and due to wind waves (index "w"), (Grant and Madsen, 1979; Signell et al., 1990).

The former was computed in the circulation model using quadratic friction law:

$$\tau_b^c = \rho C_{Db} U^2 \quad (10)$$

where U is the current velocity in the lowest water layer and C_{Db} is the friction coefficient at the bottom. Assuming a logarithmic current profile yields

$$C_{Db} = \left(\frac{1}{\kappa} (1 + \ln(z_0/\Delta h)) \right)^{-2} \quad (11)$$

where $\kappa = 0.4$ is the Von Kármán constant, Δh is the thickness of the lowest water level and $z_0 = 2.083 \times 10^{-5} m$ is the roughness length for mean grain size of $25 \mu m$ (Gayer et al., 2006, Soulsby 1997).

The contribution of wind waves to sediment dynamics uses the well-known concept that the turbulent kinetic energy associated with wave-induced bed shear stress can be described (like in the steady flows) as proportional to the second power of the orbital velocity (Jonsson, 1966):

$$\tau_b^w = 0.5 \rho f_w U_w^2 \quad (12)$$

where U_w is the maximum of the horizontal orbital wave velocity at a water depth h . Orbital velocity can be computed from the linear wave theory assuming an equivalent sine wave for the significant wave height H_s , provided peak period T and wave number k are known as:

$$U_w = \frac{\pi H_s}{T \sinh(kh)} \quad (13)$$

The wave friction factor f_w is determined by the maximum of the rough bed friction and the smooth bed friction coefficients computed as in Gayer et al.(2006).

The wave spectra properties(H_s , T , k) are computed by WAve Prediction Model WAM - a third generation non-stationary surface wave prediction model (G.J.Komen et al, 1994, Günther et al.,1991, WAMDI Group., 1988). The horizontal resolution of WAM, the used bottom topography, as well as the surface wind were the same as those used in the circulation model, so that the waves are fully consistent with the wind (see figure 6a,b and figure 7b,c). The WAM setup is essentially the same as the one used by Gayer et al. (2006). Numerical simulations agree with the dominating consensus that the Black Sea wind waves are characterised by large seasonal variability(figure 7c), the winter storms are much more frequent than the summer ones (see also Abdalla and Ozhan, 1999; and Cherneva et al. 2008). In

the open sea winds from the North and East prevails (figure 6a), generating extreme waves coming from these directions (figure 6b').

Bottom distribution of SPM

To our knowledge, there is no data available Black Sea-wide like in some other ocean basins (North Sea for instance, see Gayer et al., 2006) to specify the bed properties. A naive assumption for unlimited source of SPM at the bottom would lead to an overestimation of the resuspension for both cohesive and non-cohesive sediment (Wiberg et al., 1994, Sanford and Maa, 2001). One solution to this problem is to limit the ability for resuspension at the sea floor as a function of depth in a similar way as this has been done for the Baltic Sea by Gayer et al. (2006). Formally this approach assumes that in equation (5):

$$M_e = M_e(x, y) = M_{e,b} * \theta(x, y) \quad (14)$$

so that the "background value" $M_{e,b}$ is reduced by a topography dependent factor θ :

$$\theta(x, y) = \begin{cases} 0 & h(x, y) \leq h_{min} = 2m \\ \frac{h(x,y)-h_{min}}{h_{max}-h_{min}} & h_{min} < h(x, y) < h_{max} = 150m \\ 1 & h_{max} \leq h(x, y) \end{cases}$$

which is linearly proportional to the sea depth. The topography dependence is calibrated in a way that proposed parametrisation reduces resuspension to zero in the very shallow waters ($h < h_{min} = 2m$), below $h_{max} = 150m$ $\theta = 1$. A number of sensitivity experiments was first carried out in order to ensure reasonable fluxes consistent with the distribution of bottom sediment given by Ross et al. (1978).

The chosen value for the empirical parameter $M_{e,b} = 3.7 \times 10^{-7} kgm^{-2}s^{-1}$ is lower than other studies (Clarke and Elliot, 1998; Wang and Pinardi, 2002) but similar to Stanev et al. (2007) and shows good results.

2.3 Experiments

Experiment name	FSE	CSE	WCSE	RCSE
Sedimentation speed w_s [m/s]	10^{-5}	10^{-4}	10^{-4}	10^{-4}
Bed shear stress τ_b	$\tau_b^c + \tau_b^w$	$\tau_b^c + \tau_b^w$	$\tau_b^c + \tau_b^w$	τ_b^c
Sediment input from rivers	yes	yes	no	yes

Table 1: Summary of the performed experiments in this paper.

Several numerical experiments summarised in table 1 were performed. The first two analyse the transport of different types of sediments. The former - Fine Sediment Experiment (FSE) with settling velocity $w_s = 10^{-5}ms^{-1}$ represents very light sediment, which can be considered as a product of biological activity. The Coarse Sediment Experiment (CSE) is representative for heavier sediments delivered by rivers with typical sedimentation speed $w_s = 10^{-4}ms^{-1}$. The above sinking velocities are in the range of the once proposed in other similar studies (Wang and Pinardi, 2002, Gayer et al., 2006). More information about parameters describing sedimentation in other coastal seas are given in Mehta (1988), Clarke and Elliot (1998), Puls et al. (1999), Wang and Pinardi (2002), Mikkelsen et al. (2007), Milligan et al. (2007). To analyse the significance of river discharge and resuspension from the bottom we perform two more experiments. In one of them named Wind Coarse Sediment Experiment (WCSE) sediment input from river mouths was neglected. On the contrary, in River Coarse Sediment Experiment (RCSE) the sediment flux from rivers was enabled, but the forcing associated with the wind waves was neglected.

2.4 Satellite observations and data processing

Satellite data give a valuable source of information to study temporal and spatial characteristics of TSM. The Medium Resolution Imaging Spectrometer (MERIS) which operates on the environmental remote sensing satellite ENVISAT is used to retrieve the concentrations of phytoplankton pigment, suspended matter and yellow substances ("gelbstoff") from water leaving radiance reflectance (Doerffer and Schiller, 1998; Doerffer et al., 2006). The used algorithm is a multiple non-linear regression method ("Neural Network") that needs as an input the directional water

336 leaving radiance reflectances of 8 MERIS bands and solar and observation angles.
337 Its coefficients are determined using a feed forward back-propagation optimisation
338 ("training") technique. Two key features of the presented algorithm justified its
339 choice: (1) it treats a wide range of concentrations which cover open ocean case I
340 water as well as turbid case II waters, and (2) it is capable to separate the suspended
341 matter from the other substances. The Basic Toolbox for ENVISAT (A)ATSR and
342 MERIS (BEAM)(Brockmann, 2003), provided by ESA ENVISAT project includes
343 plugin with the latest validated version of this Neural Network. We had applied it
344 over the raw MERIS level 1 data for 2007-2008 period to produce TSM concentration
345 in the Black Sea (figure 11) that is to produce MERIS level 2 data.

346 The propagation of turbid waters and their variability in time and space was inti-
347 mately related to sediment transport and remobilisation. Daily mapping of satellite
348 data can be used to derive important statistics, observe the extension of river plumes,
349 and the re-suspension of sediments from the bottom in shallow waters, usually due
350 to wind-wave action. The data is available as a MERIS level 3 product generated
351 after re-processing the MERIS level 2 data. The detection of turbid water is based
352 on a threshold value of water reflectivity at a given wavelength of the spectrum
353 equivalent to a sediment concentration of 0.75g m^{-3} (Moore et al., 1999) and rep-
354 represents an operational product of the Coastal GMES Service Element Coastwatch.
355 Black Sea Coastwatch data were further used here for inter-comparisons of model
356 versus observations (see further in text).

3 Results

3.1 Horizontal circulation

359 The realism of the numerical simulation of circulation has been analysed in details
360 by Grayek et al. (2010) thus only a brief overview is presented in this paper. The
361 simulated currents (figure 8) agree with previous simulations based on different mod-
362 els (Staneva et al., 2001; Stanev et al., 2003). They reveal a dominating cyclonic
363 pattern (figure 8a) that weakens during spring and summer (figure 8c, d). The lat-
364 ter is exemplified in figure 8c, d by the change of the sign of anomaly-circulation
365 (clockwise). The anomaly displays cyclonic sense of rotation in winter (figure 8b),

366 that is the Rim Current gets stronger. The range of the winter-to summer change
367 in the intensity of circulation is about the half of the annual mean value. The pre-
368 liminary expectation is that this type of circulation will support a counter-clockwise
369 transport of SPM.

370 3.2 Bed shear stress

371 Sediment dynamics was dominated by the distribution of sources and transport
372 processes. The bottom acted as a source or sink depending on whether τ_b was larger
373 or smaller than τ_{crit} . Therefore, it is instructive to consider the temporal and spatial
374 characteristics of τ_b . The annual mean patterns of τ_b^c and τ_b^w demonstrated that both
375 constituents of the bed shear stress have maxima in the coastal area (figure 9a,b).
376 Over most of the bottom area that is in the deep ocean, the bed shear stress was
377 more than one order of magnitude smaller than the averaged shelf value. Large
378 values of τ_b^c along the western and southern coast were found in the embayment
379 and headlands (e. g. Cape Kaliakra), while bed stress along the Caucasian coast
380 was estimated to be extremely small, which is explained by the steep bottom and
381 absence of pronounced headlands and embayments. Overall, figure 9c gives an idea
382 about areas where the annual mean $\tau_b > \tau_{crit} = 0.02N/m^2$, that is areas favourable
383 for erosion.

384 The temporal variability of the bed shear stress induced by wind waves τ_b^w (not
385 shown here) follows the temporal variability of the wind-wave forcing (figure 7c),
386 larger values are computed in winter and smaller ones in summer. It is important to
387 mention that the action of the wind waves is relatively shallow and decreases rapidly
388 with increasing ocean depth. The bed shear stress due to currents τ_b^c has a distinct
389 seasonal cycle. During the intensification of the Rim Current in the autumn and
390 winter τ_b^c reached the critical values all along the western coast.

391 More detailed analysis showed that there is no clear correlation between wind forcing
392 (not the wind-wave one) and bottom friction. It appears that the increased values
393 of bottom friction due to currents are a result of the eddy variability. The episodes
394 of extreme τ_b^c that drive sediment erosion were explained by the events in the ocean
395 weather, not the atmospheric weather.

3.3 Temporal variability

Temporal variability simulated in CSE and FSE is displayed in figure 10 for area mean concentrations and concentrations on the shelf and in the open sea (see figure 1 for the positions). In both runs area mean concentration was dominated by the high concentrations in the near coastal locations, therefore the black curve is above the green one. Coarse SPM reached lower concentrations because much of it deposited rapidly on the bottom.

In both runs maximum concentrations were observed during winter months, identifying the clear response of sediment system to atmospheric forcing. The signal was much smoother for the fine sediment (compare figure 10a and figure 10b), which is explained by the fact that small sinking velocities keep SPM longer in the suspension. In the case of coarse sediment shallow depths and large sinking velocities led to a removal of SPM from the suspension very rapidly therefore the oscillations were more pronounced and directly linked to the atmospheric forcing. Oscillations tend to decrease in magnitude with the increasing depth (compare the green and black lines in figure 10a).

The shape of the SPM curves in the sampled locations (figure 10c,d,e,f) differs significantly from the area mean ones - the deeper layers in most cases displayed also larger concentrations as compared to the sea surface. This is consistent with the fact that sediment settled and propagated mainly by advection. We have to denote that the surface maxima in the shelf point (figure 10c) is caused by direct transport from the Danube river because of particular circulation conditions. At the same time higher concentrations at 50m in the open sea were not due to local sources (if any they are deep below sea surface).

One interesting phenomenon displayed by the numerical simulations deserves to be mentioned. There were short time intervals in January, February and December, during which the two sediment curves overlapped. This could be considered as an indication of the role of convection homogenising rapidly the water column (figure 10c,e). This could have shaped SPM stratification and transport. This situation seems very "exotic" because, if the examined process occurs in realistic systems it has to work against settling, that is mixing associated with convective instability, could bring SPM up to the sea surface.

One last conclusion from figure 10f needs also to be made: in the open ocean area,

concentration of coarse SPM was negligible, that is this fraction sunk so rapidly in the coastal zone that it could not be found in the surface layers of ocean interior. On the contrary, the light sediment transported by advection is still in suspension in the open sea (figure 10e).

3.4 Horizontal sediment patterns. Comparison between simulations and observations

3.4.1 Observations

In the following we describe six scenes for 2007, which we sampled from reliable data of more than 400 relatively cloud-free images. The sampling criterion included as much as possible cover, larger diversity in the spatial characteristics of meso-scale features, clear agreement or disagreement with numerical simulations. For convenience the data from Azov sea was masked because of extremely high concentrations of SPM from biogenic origin.

The major source of sediment was observed in the area of the north western shelf (figure 11a, c, e, f). The region with largest values was in a close proximity of the Danube Delta. From there the sediment plume propagated in most cases to the south along the Romanian and Bulgarian coasts and tended to form small-scale “fingers” (figure 11c, f) about 10-20 km long and 2 km wide. These “fingers” identified the area of origin of high concentration, which extended in some cases far into the open ocean.

In some periods, as for instance during 21 July (figure 11e) the Danube plume turned to the North and high TSM values dominated the entire shelf area. In this case a very thin intrusion formed at the western coast of Crimea, penetrating very deep into the basin interior. Similar propagation feature (almost on the continental slope) was clearly seen in figure 11f separated from the rest of shelf water by a zone with lower concentration.

The second and third important areas of origin of TSM are south of the Kerch Strait (figure 11b) and North of the Turkish Coast (Cape Synop, figure 11c). As illustrated by the model, in these areas concentration increases as a consequences of a combination of atmospheric forcing and dominating meso-scale dynamics in the ocean. As seen in figure 11c, a chain of narrow intrusions propagated to the west

460 between the Rim Current and the coast.
461 TSM patterns allow tracing important elements of the Black Sea eddy/meso-scale
462 variability. Larger eddies had scales of several tenths of kilometre (figure 11b) and
463 the corresponding mushroom structures and intrusions reached scales of more than
464 100 km. The size of coastal waves on the Rim Current and eddies propagating
465 between the Rim Current and the coast was clearly detectable as about 50-80 km
466 (figure 11b, c). In some cases, as for instance in figure 11f headlands along the
467 Asia Minor coast appeared to give the origin of very small scale scale eddies with
468 increased TSM concentration.

469 3.4.2 Numerical simulations

470 Because we do not assimilate data we do not expect that the individual synoptic
471 images from the model agree with observations that is the simulated eddies were not
472 positioned at the same place and time as the observed ones. More important here is
473 to discuss general characteristics of TSM, such as scales, types of patterns, sources
474 and their impact on the spatial characteristics. Therefore we show in figure 12
475 several snapshots from the model, which illustrate its performance, in particular
476 characteristic features comparable with the ones derived from the observations. We
477 also admit that distinguishing between sediment of different origin (river, bottom,
478 biology) is not easy, which could create problems in the inter-comparison.

479 The overall comparison between simulations and observations indicated that the
480 model was still too coarse to resolve the small scales of 1-3 km, which dominated
481 the near coastal dynamics of sediment. These features discussed above were sub-grid
482 for our model.

483 Eddies formed coherent flow structures with dimensions similar to coastal embay-
484 ments. They propagated through these embayments over time scale of days to weeks
485 enabling the transport of suspended sediments over long distances. The consecutive
486 snapshots from the model gave an overall idea of the magnitude of changes that
487 could occur over only a few days.

488 There was a good agreement between observations and simulations, as this concerned
489 the north-western shelf. This is the area where concentrations reached in the model
490 local maximum at all times. The expected maximum in the Danube mouth observed
491 by the satellite (figure 11a,c,f) was simulated also by the model(figure 12b,c,d,f).

492 Large-scale eddies, intrusions, mushroom-type patterns (figure 12a, c,d,f) resembled
493 the observations (compare with figure 11e,f). Comparison between figure 12a and
494 figure 11a could serve as an indication that the model also replicated the intrusions
495 associated in the area east of Kerch Strait.

496 Figure 12f gives an idea of the evolution of coastal waves and formation of coastal
497 eddies on the continental slope. Similar process has been earlier observed in the same
498 area by Sur et al. (1996) in their analysis of satellite data. In the present model
499 validation figure 12f gave the model analogue of the chain of eddies in figure 11c
500 propagating in the area of the western Black Sea continental slope.

501 From figure 12c, d it becomes clear that the sediment source along the Asia Minor
502 propagated in a similar way as in figure 11b,c,e, however fine scale zonal intrusions
503 could not be resolved by the model. Simulated coastal eddies along Asia Minor coast
504 similar to ones in figure 11b, c were however resolved in figure 12b, e. An eastward
505 propagation of small scale coastal waves between Sakarya and Sinop (figure 12d)
506 found also a support in the satellite data analysed in the past (Sur et al., 1996).
507 A rapid sinking of the sediment coming from the rivers along the Asia Minor coast
508 in the vicinity of the vertical continental slope, which was identified by the abrupt
509 decrease of SPM (figure 11f) was also visible in the model simulations(figure 12c).
510 However, a number of small scale features are missing in the simulations indicating
511 the need to further increase model resolution.

512 It is noteworthy that depending on the local circulation in the easternmost part of
513 the Black Sea, which is sometimes dominated by the so called Batumi eddy, the
514 source of sediment from the Georgian Rivers propagated either to the north west
515 (along the Caucasus coast, figure 12f) or to the south (figure 12a).

516 **3.5 Sediment transport**

517 **3.5.1 Penetration of SPM into the basin interior. Vertical sediment** 518 **patterns.**

519 Sediment modelling provides a means of illustrating how matter from rivers or remo-
520 bilized from the bottom may be redistributed. A number of studies (e. g. Lewis and
521 Landing, 1991) claim that the intrusions from bottom represented a mechanism by
522 which bottom water penetrates the basin interior (see also Konovalov et al., 2003).

523 So far, this issue has not been thoroughly addressed by the Black Sea numerical
524 modelling. Our simulations clearly revealed traces of such intrusions (figure 13) by
525 meso-scale eddies along the steep bottom.

526 Obviously, the model revealed the key mechanisms for transport and delivery of
527 sediment from the shelf environments into the deep ocean. It mimicked aspects of
528 long-term depositional pattern observed in Holocene stratigraphy (see Fig. 7 from
529 Ross et al., 1978). Danube sediments travelled southward during a series of episodic
530 transport events. However, their penetration into the basin interior appeared to be
531 limited.

532 Figure 13 displays annual mean vertical sections of SPM along 31 E simulated in
533 FSE and CSE after 3 years of integration (trends in the sediment model are not
534 seen after 2 years of integration). The difference between these plots explains the
535 role of settling velocity for the penetration of SPM in the basin. In these two
536 experiments the critical bed shear stress was always reached somewhere, thus the
537 sediment sources were associated both with fluxes from rivers and remobilisation of
538 bottom sediment under extreme events (providing extreme bed shear stress). For a
539 relatively large settling velocity (CSE) the sediment plume near the ocean surface
540 could not propagate far from the river mouths and beyond the continental slope.
541 This means that the rapidly sinking sediment could only leave the coastal area in
542 the deeper layers (figure 13b). This is indeed the case in the southern and eastern
543 Black Sea, which lacks a continental shelf, therefore sediment from coastal area could
544 easily reach larger depths.

545 Differently from CSE, in FSE the coastal sediment laden water propagated on the
546 shelf much farther into the direction of the continental slope. The specific balance
547 between the horizontal transport from one side and settling from another enabled
548 coastal intrusions of suspended sediment to reach the basin interior (compare fig-
549 ure 13a with figure 13b) building a layer with maximum concentrations at ~ 200 -300
550 m. The position of this maximum is explained by the fact that most of the coastal
551 intrusions appear in the area of continental slope at about 200 m, that is where the
552 Rim Current is usually observed. This finding is supported by the observations of
553 Wakeham et al. (1991), as well as by the recent unpublished data collected in the
554 frame of SESAME project ([refhttp://www.sesame-ip.eu/](http://www.sesame-ip.eu/)).

555 In the Black Sea the steepness of the continental slope is extremely large, in particu-

lar in its southern and eastern parts, and the topographic control (JEBAR, Stanev, 1990), which tends to trap the jet over the continental slope weakens. As a consequence of this specificity (vertical wall) the instabilities of the Rim Current increase in magnitude (Stanev, 2005). As demonstrated in section 3.4 both observations and numerical simulations revealed long-reaching horizontal intrusions of coastal waters in the surface and intermediate layers caused by jets and coherent eddies, which played the major role in the propagation of SPM further into the open ocean.

3.5.2 Erosion and Deposition

The above discussion is extended below with the analysis of erosion and deposition of SPM (figure 14). As mentioned in the Introduction, existing data are not sufficient to reveal the spatial characteristics of deposition, therefore numerical simulations could shed more light on this issue.

The total bed shear stress τ_b (figure 9c) governing the process of resuspension was determining the horizontal pattern of the erosion rate (figure 14a) thus a maximum was found in the entire northwestern shelf and along the western coast. The low level of turbulence in the deep ocean disabled erosion below 160m.

Substantial differences in the deposition rates are identified in the various simulations. In CSE the coarse sediment was trapped near the mouths of the rivers and in the southwestern coastal zone. On the contrary, the fine sediment (figure 14b) had the ability to be transported far from the area of its origin and significant part of the Danube input settled down in the entire northwestern shelf and along the western coast.

3.6 River fluxes versus erosion from bottom

In the analyses done so far, no separation of responses to different forcing has been addressed. In order to get an idea of the relevant importance of river discharge and seabed resuspension, we carried out the following two experiments: (1) WCSE where the sediment fluxes from river mouths were disabled (figure 16a, c, e, g) and (2) RCSE where fluxes from rivers were enabled, but the forcing associated with the wind waves was disabled (figure 16b, d, f, h). The joint action of the two mechanisms was included in the CSE and the corresponding results are shown in figure 15b, d,

f, h, along with with a similar presentation of observations from satellites.

In the absence of the wind-wave forcing, the sediment remained very close to the river mouths, and the seasonal changes were less important. Little exception can be seen during the cold part of the year when the strong currents increase the bottom shear stress and erode sediment along the western coast.

Simulations in which the wind-wave forcing was enabled but sediment flux from rivers was disabled (figure 16a, c, e, g) were very similar to CSE in the fall and winter seasons, revealing again that the sediment patterns are highly dependent upon the atmospheric forcing and resulting bottom resuspension. However, the river runoff maintains the availability of sediments, which are transported from the shelf into the open ocean.

In this context the study of Wijsman et al. (1999) on the sediment characteristics and benthic activity on the northwestern Black Sea shelf is worth mentioning. Using in-situ observations with a focus on the the River Danube as a source of suspended solids to the continental shelf, these authors identified the following regions: (1) the area just in front of the Danube Delta (large discharge), (2) the northern part of the continental shelf where the majority of the Danube discharges are transported, and (3) the southern part of the continental shelf characterised by low sedimentation rates and low rates of benthic mineralisation. Analysis provided above supports this schematics, in particular as it concerns the area of low sedimentation: in the case of absence of resuspension (figure 16b, d, f, h) sediment concentrations are high only in front of the Danube Delta. South of the delta (along the Bulgarian coast) concentration of coarse sediment decreases substantially.

4 Conclusions

The modelling examples and data from remote sensing discussed in this paper demonstrated clearly the potential for the model to address the sediment dynamics in the Black Sea. We found that forcing due to wind waves provides the major driving element, which results from episodic high energy events contributing to the increase of horizontal sediment fluxes over the north-western shelf. Circulation and eddies shaped further the distribution of sediment. Comparison with observations was encouraging for the future use of this model.

617 The penetration of suspended sediment into the basin interior was shaped by the
618 dynamics of coastal and open-ocean eddies. Fine sediment at sea surface could
619 efficiently cross the continental slope and was observed at sea surface both in open-
620 ocean observations and simulations in the form of coherent structures intrusions and
621 wave-like patterns. On the opposite, coarser fractions sink rapidly along the bottom
622 and can be observed over the continental slope and only below sea surface.

623 The patterns characterizing coarse and fine sediment are too different in the vertical
624 plane. While the coarse sediment follows approximately the bottom, the fine one
625 enters the deep ocean at about 200 m building a pronounced layer of maximum
626 concentration in the basin interior. The plausibility of numerical simulations and
627 interpretation of this result is supported by earlier local observations (Wakeham et
628 al., 1991), as well as by recent surveys in the frame of SESAME activities.

629 Some unexpected results deserve also to be mentioned. It appeared that winter con-
630 vection mixes SPM in the surface layer, and could provide an important mechanism
631 for sediment transport. This mechanism identified by the numerical simulations
632 needs to be validated against dedicated observations.

633 Numerical experiments with included or excluded different forcing terms, e. g.
634 sediment flux from rivers appeared to be very useful enabling determine the spatial
635 extensions of different point-sources. Analysis of numerical simulations supported
636 earlier concepts of Wijsman et al. (1999) about sedimentation around of and south
637 of the Danube Delta. Without resuspension mechanism in the model, sediment
638 concentrations were high only in front of the Danube Delta. The continental shelf
639 south of the delta (along the Bulgarian coast) was characterised by low concentration
640 of coarse sediment and low rates of sedimentation.

641 Some future developments have to be performed in order to construct a realistic map
642 of bottom sediment, which is needed to improve the bottom boundary condition.

643 Once a map is available combining different type of sediments in a common system
644 could provide more accurate information about the erosion process. Coupling the
645 sediment transport model to ecosystem model (e. g. Gregoire and Beckers, 2004)
646 will increase the realism and open the road to studying the light climate in this
647 basin. Testing the model simulations against ship-based and moored instruments
648 measurements of currents and sediment transport is also needed. Downscaling to
649 scales on the order of 100 meters to permit adequate resolution of bottom topography

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650 and complex shoreline has also to be tried.

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5 Figures

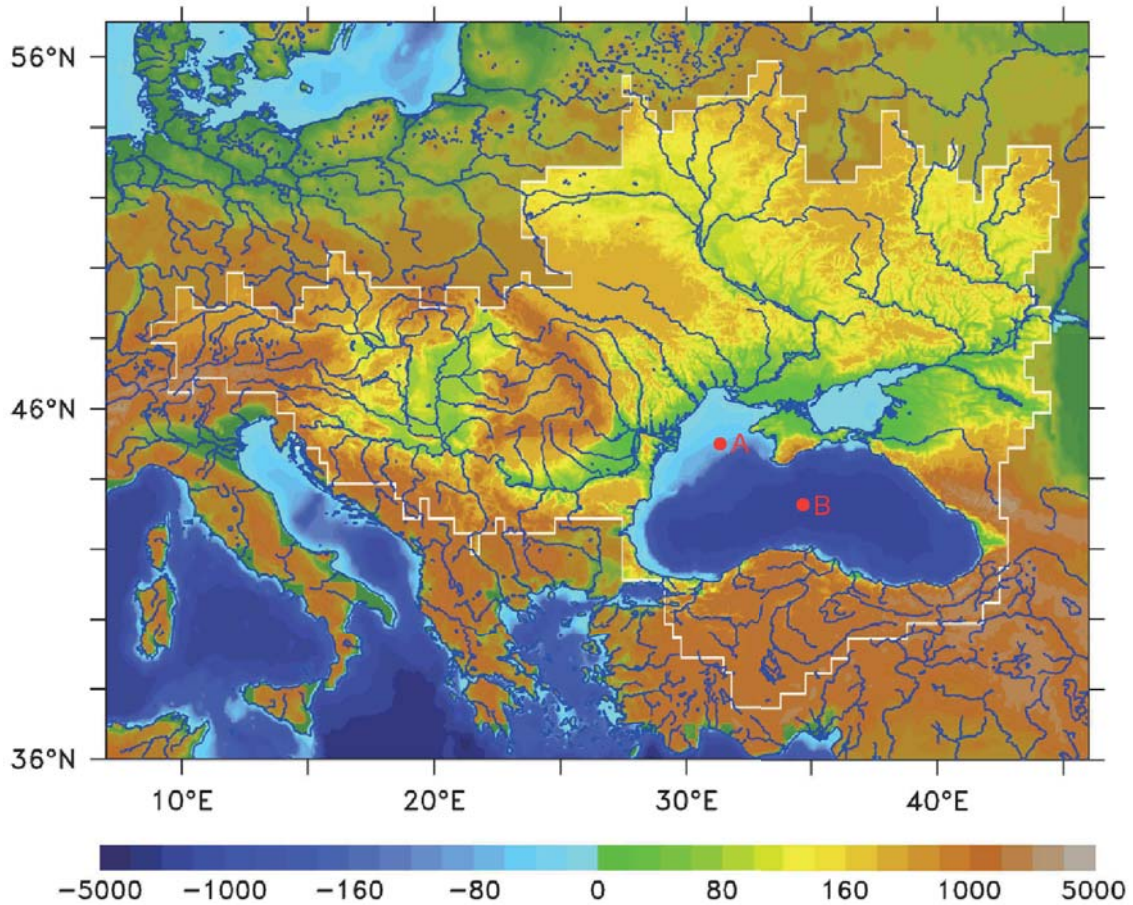


Figure 1: The Black Sea and its catchment area. Orography of the catchment area (m) is plotted with brighter colours. Bottom topography (m) is coloured in blue, shelf in light blue, deep-ocean in dark blue. Note that shelf is almost absent in the southern and eastern parts of the sea. Positions of locations discussed further in the paper are also shown.

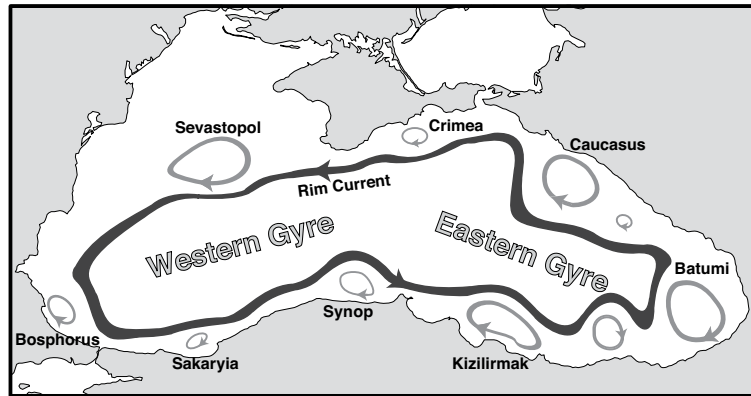


Figure 2: Black Sea circulation (presented schematically based on numerical simulations of Stanev, 2005). The names of major coastal eddies as known from observations (usually, associated with geographic names, for example, capes and towns) are also given.

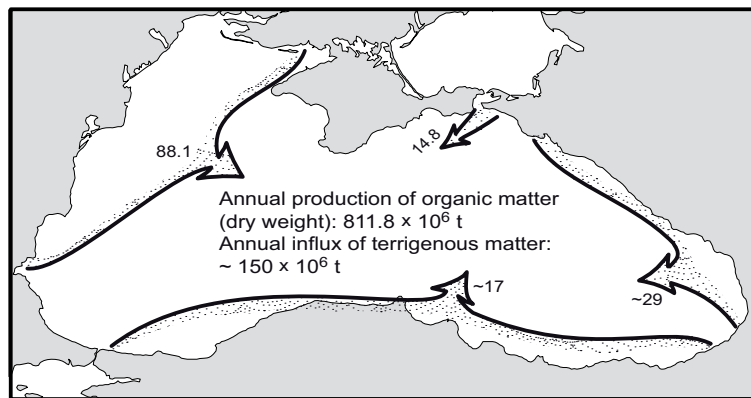


Figure 3: Supply of sedimentary material (in 10^6 t/yr) into the Black Sea interior (replotted from Shimkus and Trimonis, 1974).

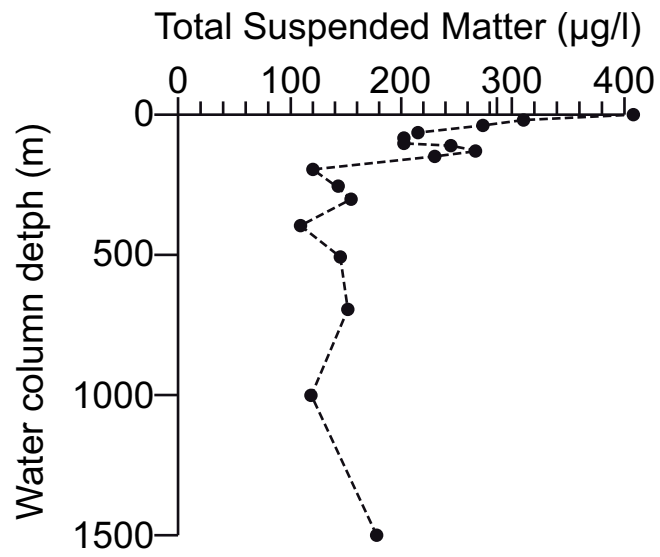


Figure 4: Vertical distribution of TSM (redrawn from Wakeham et al., 1991).

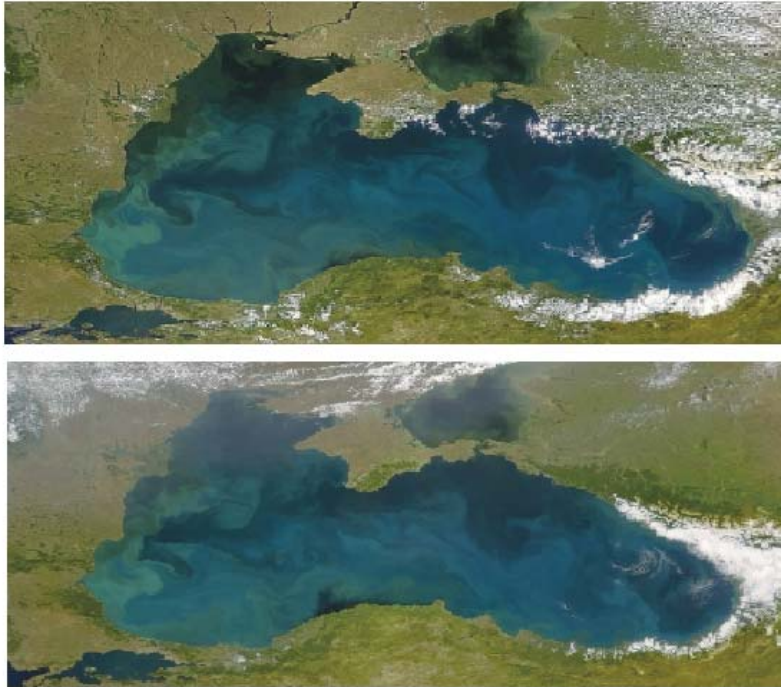


Figure 5: This image has been provided by SeaWiFS Project, NASA/Goddard Space Flight Center and ORBIMAGE and shows the surface patterns of the Black Sea turbid waters on June 11 (upper panel) and June 13 2000 (bottom panel).

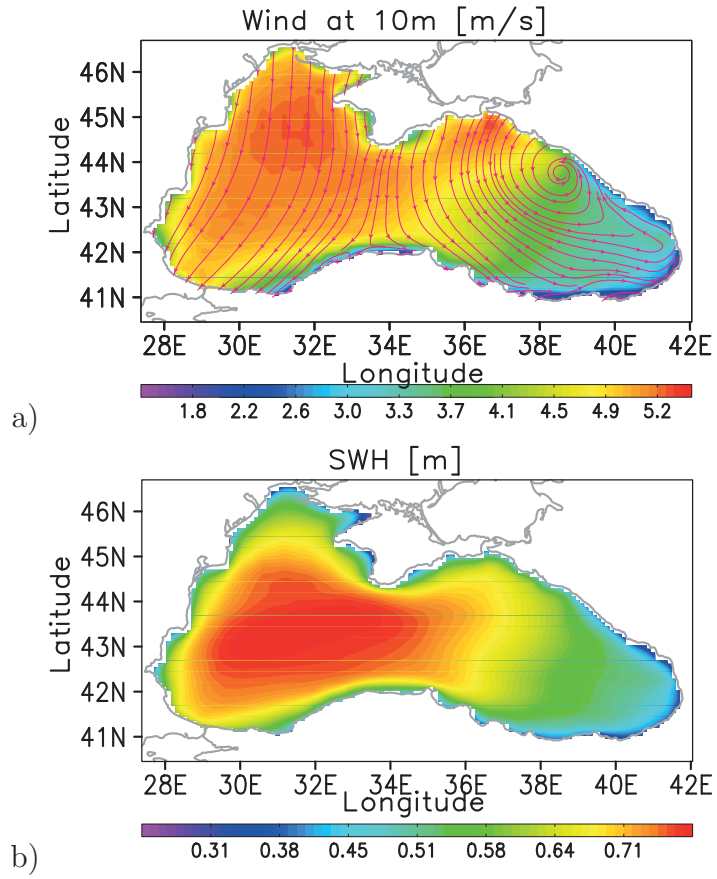


Figure 6: Annual mean wind (a) and significant wave height(SWH) forcing (b).

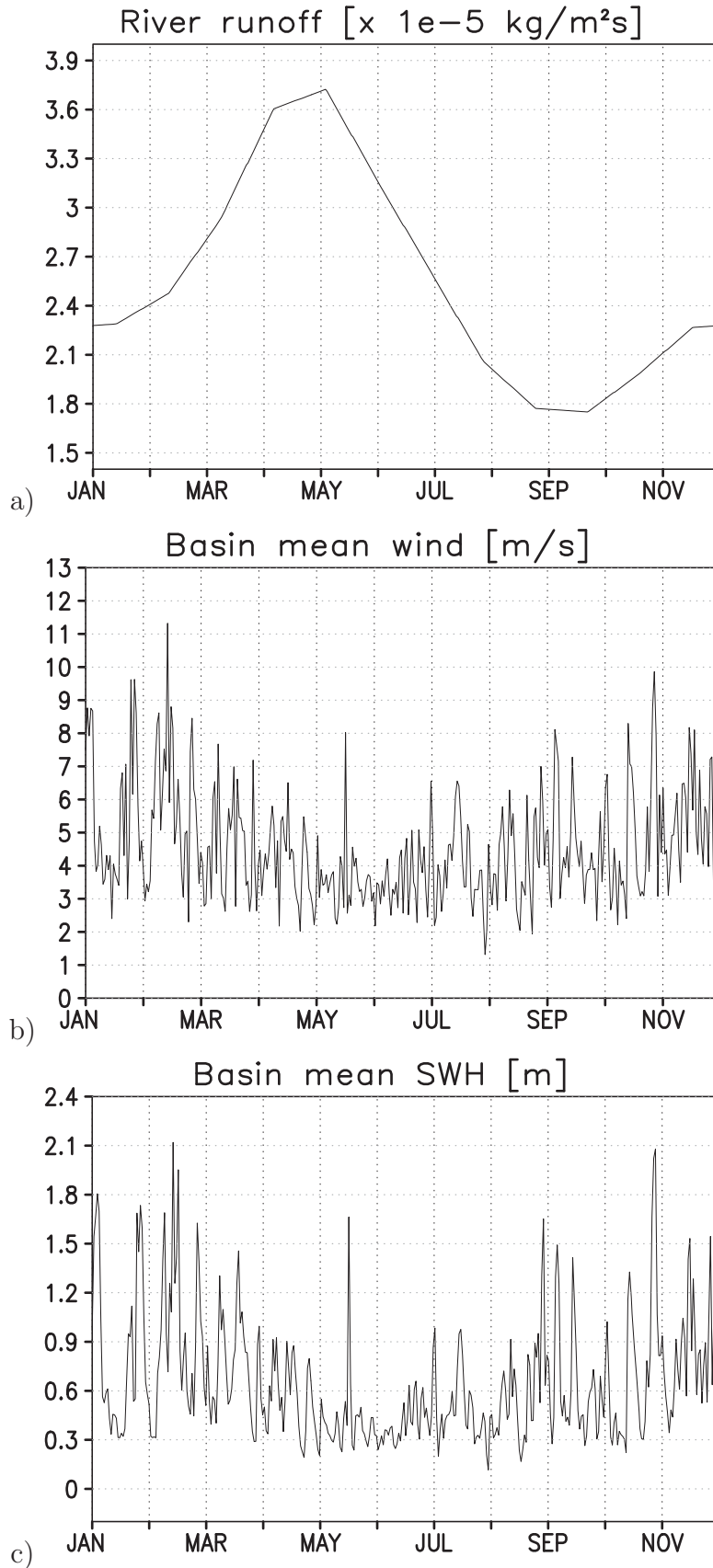
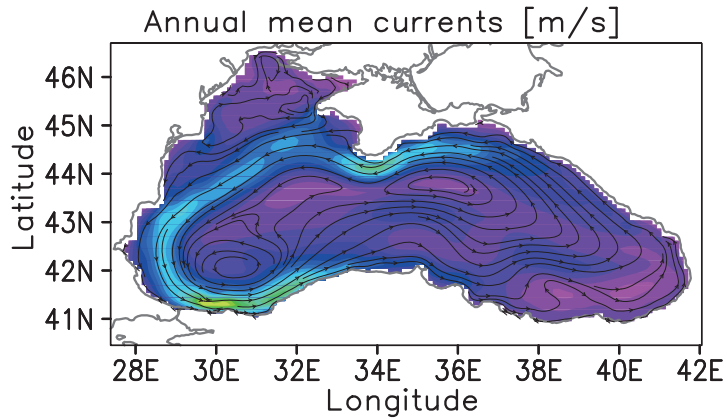
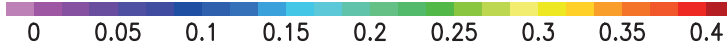


Figure 7: Temporal variability of the river runoff (a), basin mean wind (b), and basin mean significant wave height (SWH) (c).

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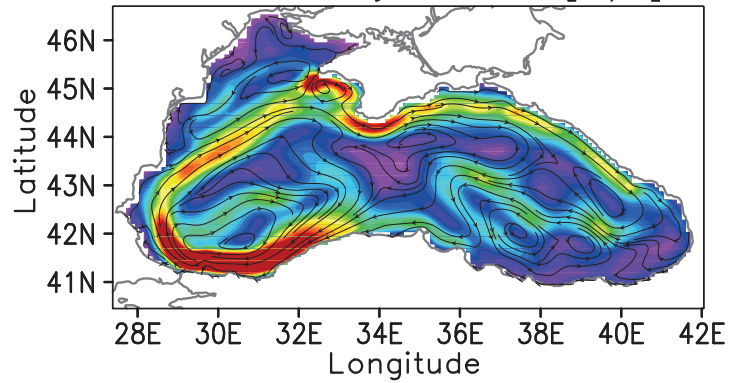
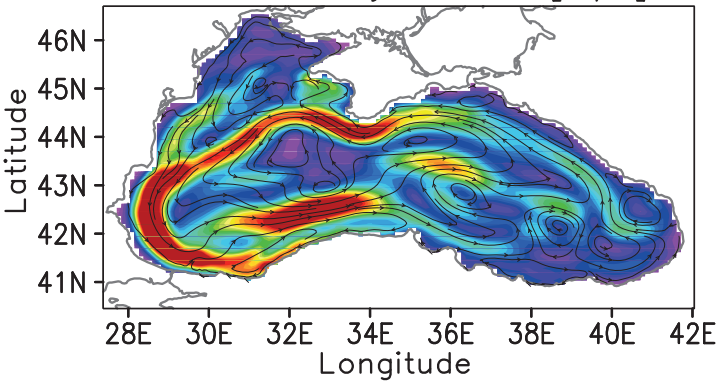


a)



Currents anomaly DJF-AM [m/s]

Currents anomaly MAM-AM [m/s]



Currents anomaly JJA-AM [m/s]

Currents anomaly SON-AM [m/s]

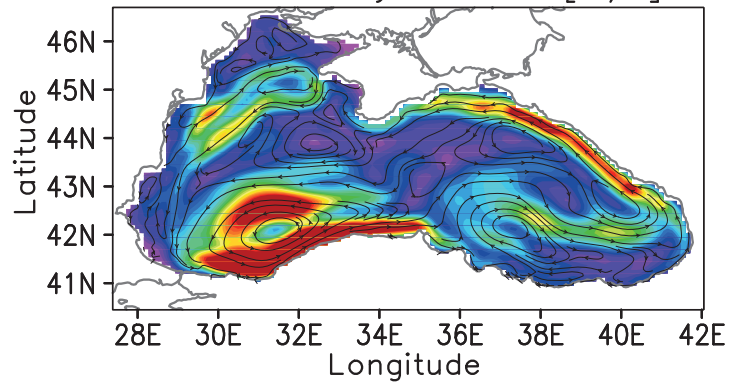
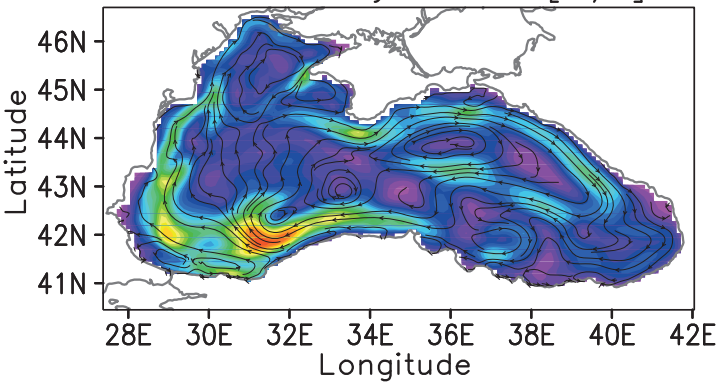


Figure 8: Simulated annual mean (a) and anomalies of the seasonal mean (b-e) surface currents.

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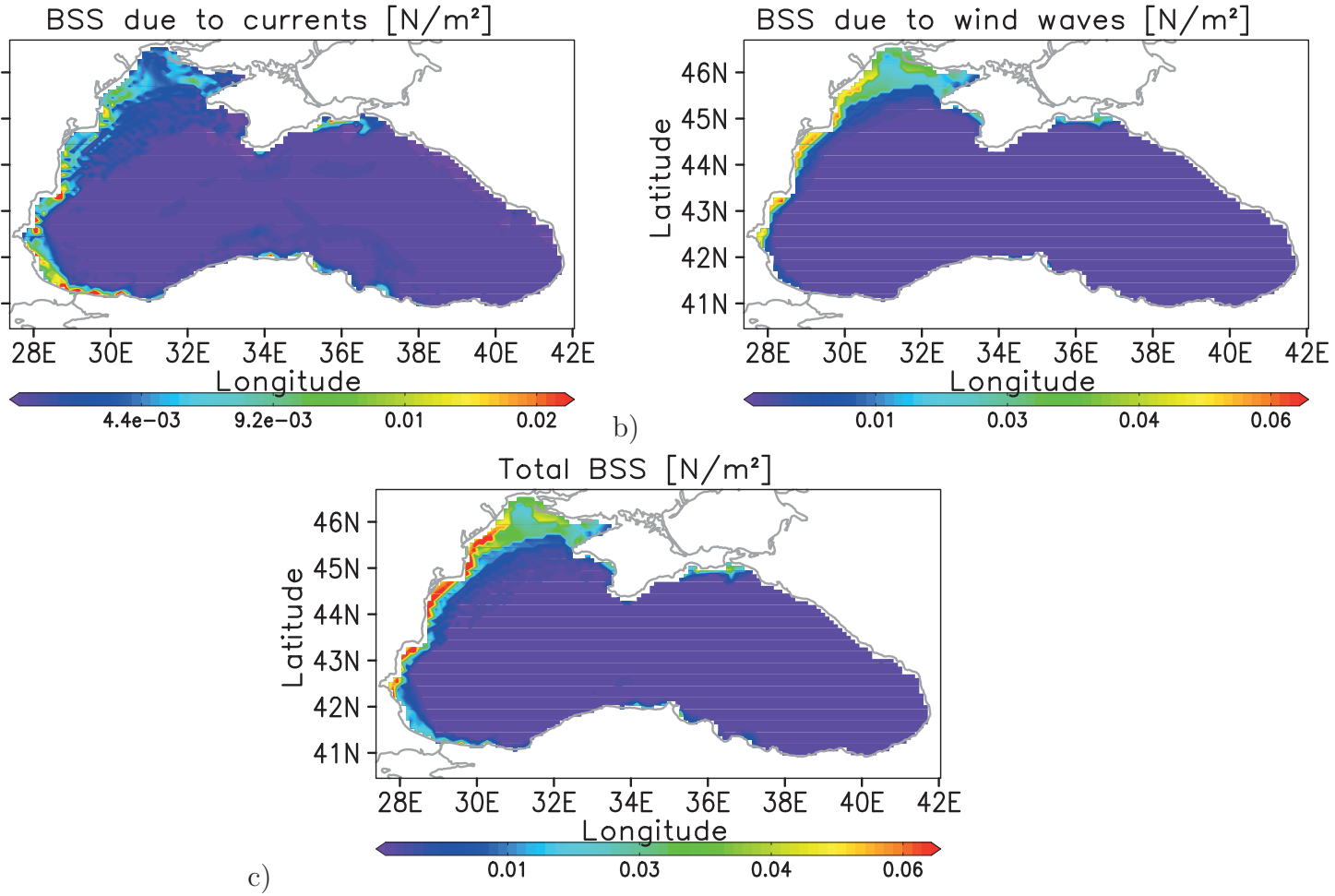


Figure 9: Annual mean bed shear stress (BSS) due to currents (a), wind waves (b) and combined (c)

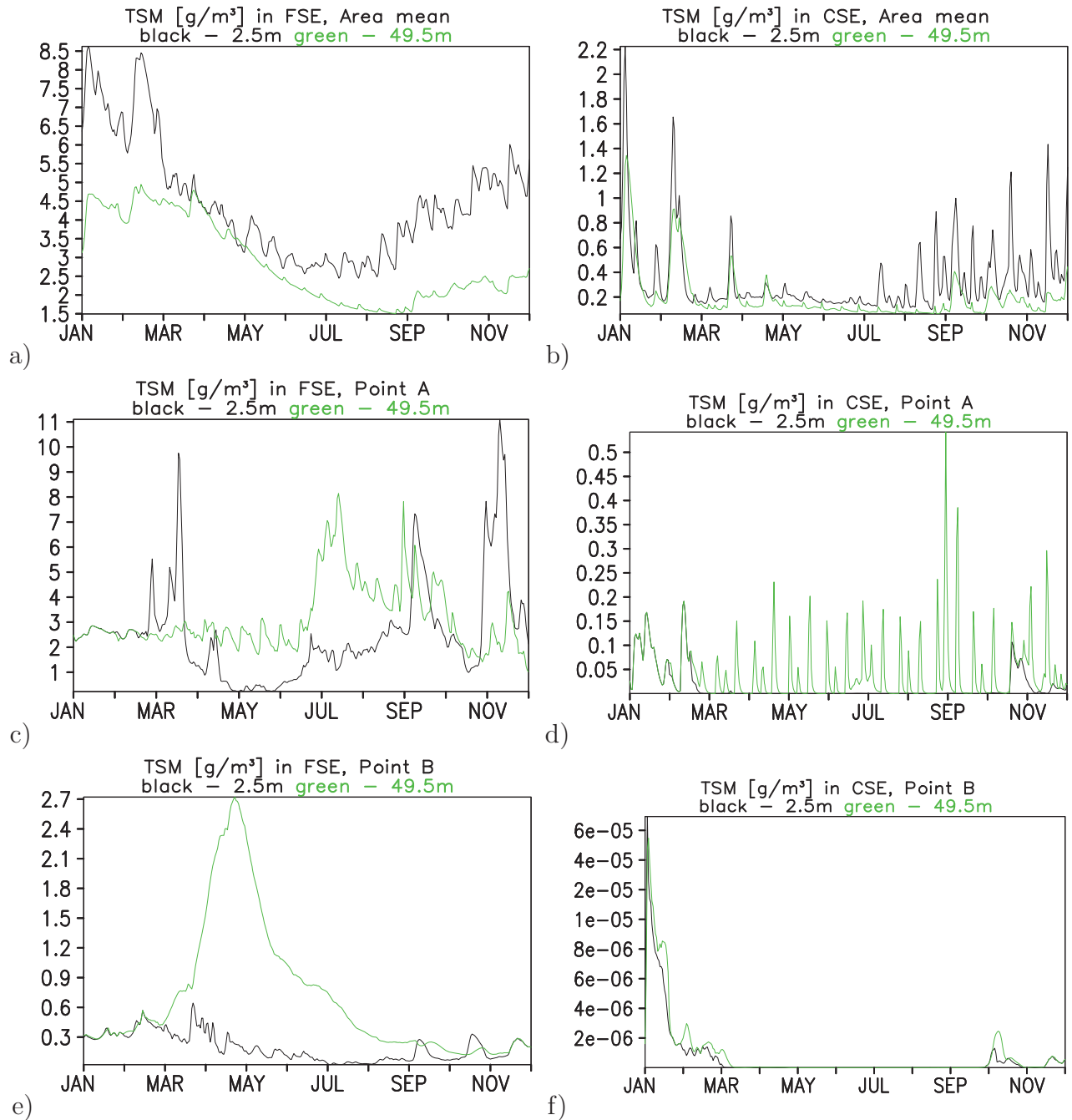


Figure 10: Temporal evolution of surface TSM (black line) and TSM at 49.5 m (green line) in FSE (a, b, c) and CSE (d, e, f). Basin mean concentration (a,d) as well as local concentration in the area of north-western shelf (Point A, 45.0N,31.0E, b, e) and open sea (Point B, 43.3N, 34.3E, c, f) are presented. For the position of the locations see figure 1.

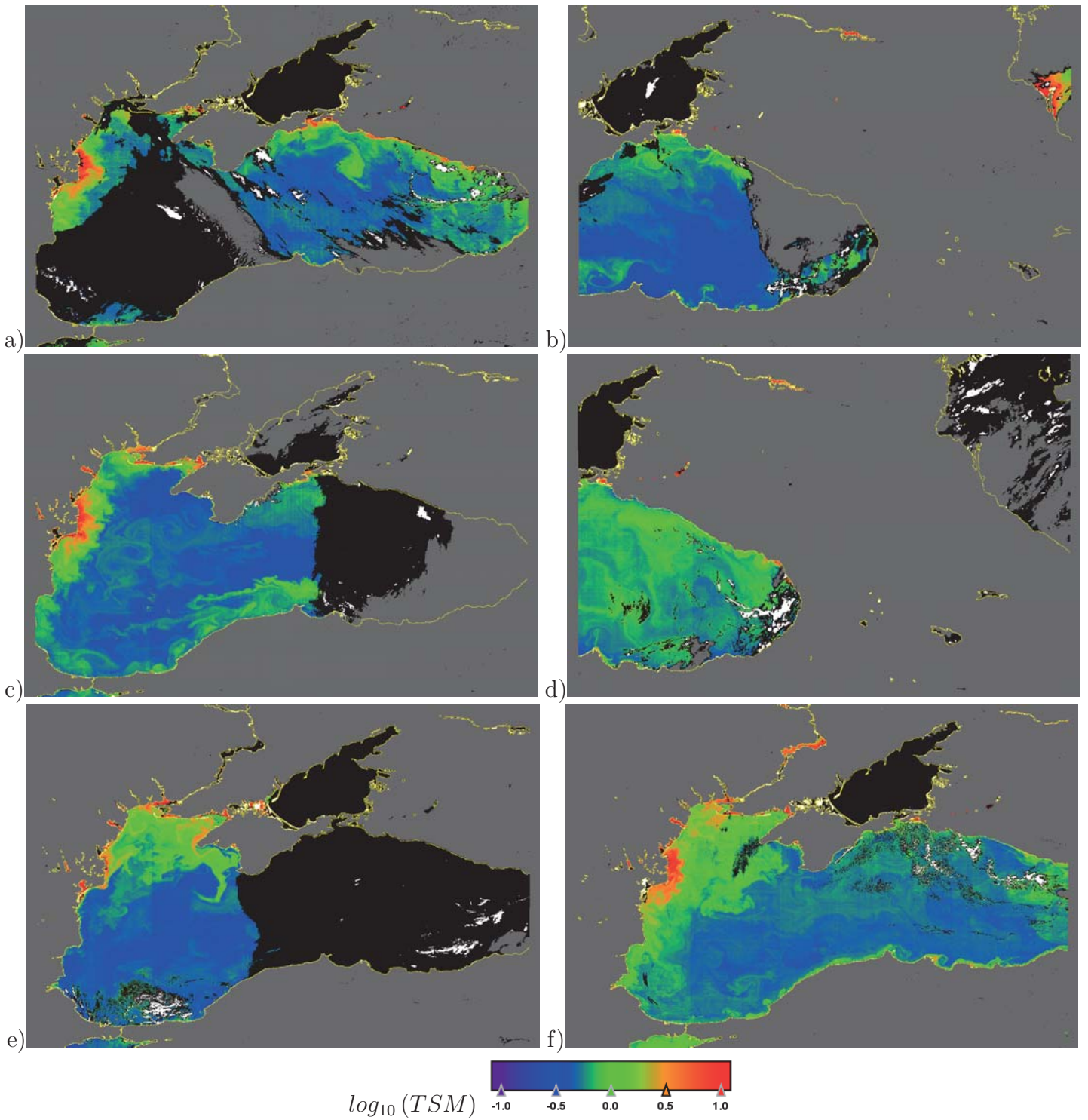


Figure 11: Surface TSM concentration from MERIS satellite in 2007 on 19.Mar. - a, 11.Apr. - b, 23.Apr. - c, 03.Jul. - d, 21.Jul. - e, 18.Oct. - f

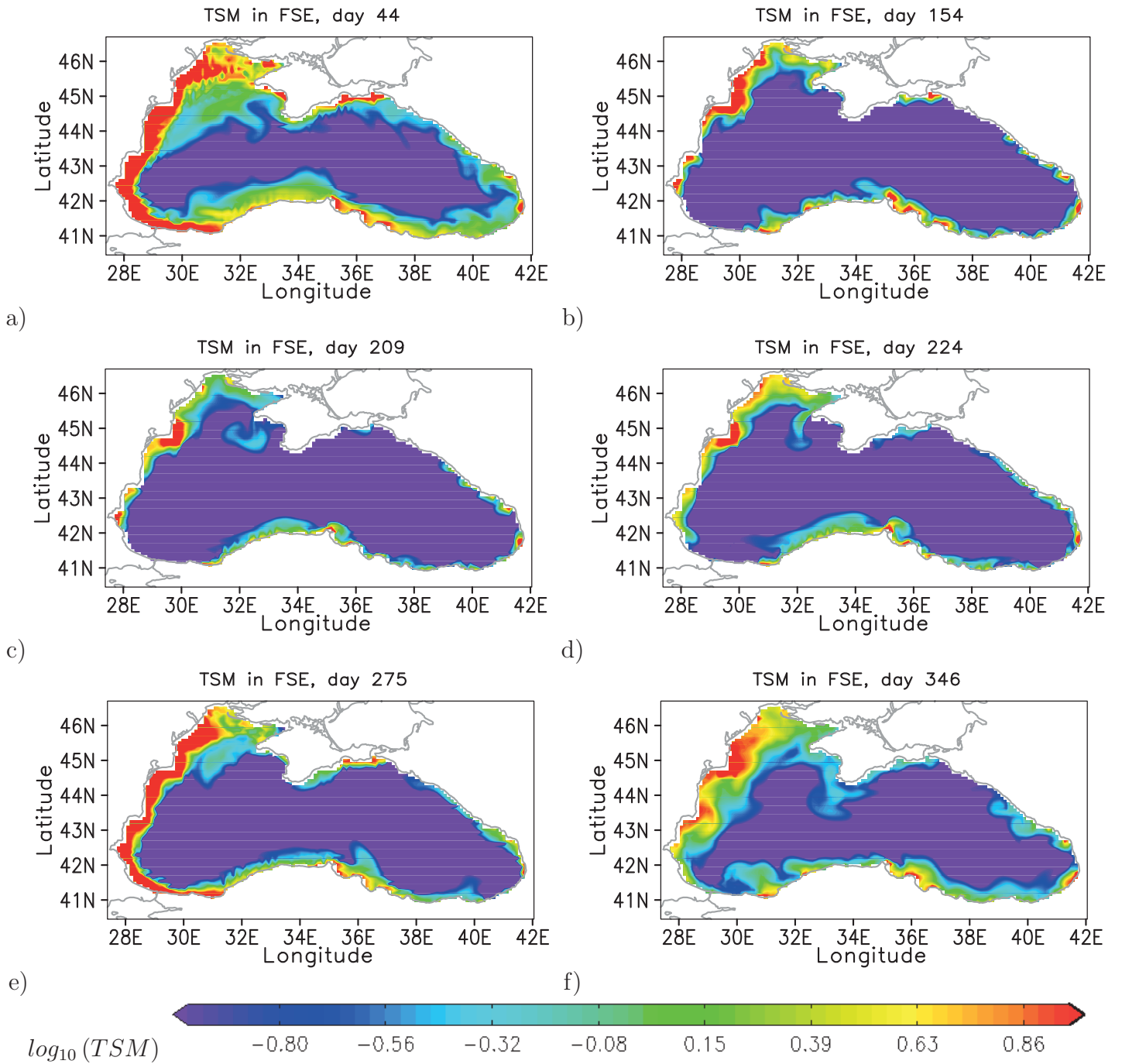


Figure 12: Surface TSM concentration as simulated in FSE in 2008 on 13.Feb. - a, 02.Jun. - b, 27.Jul. - c, 11.Aug. - d, 01.Oct. -e, 11.Dec. - f

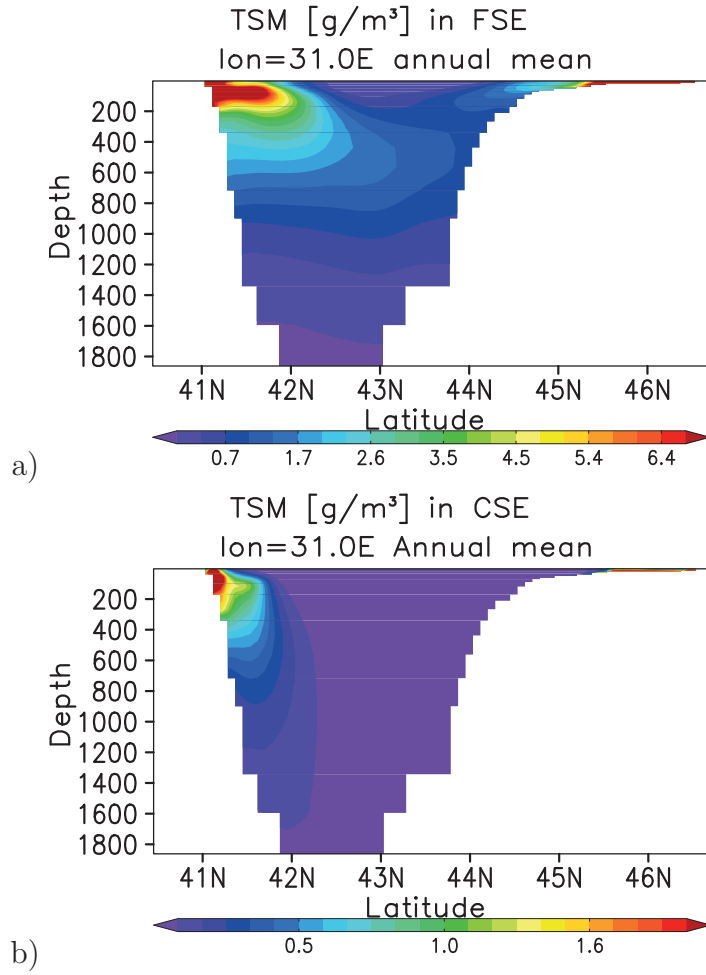


Figure 13: Vertical cross-sections of TSM at 31.0E as simulated in FSE - a) and CSE - b).

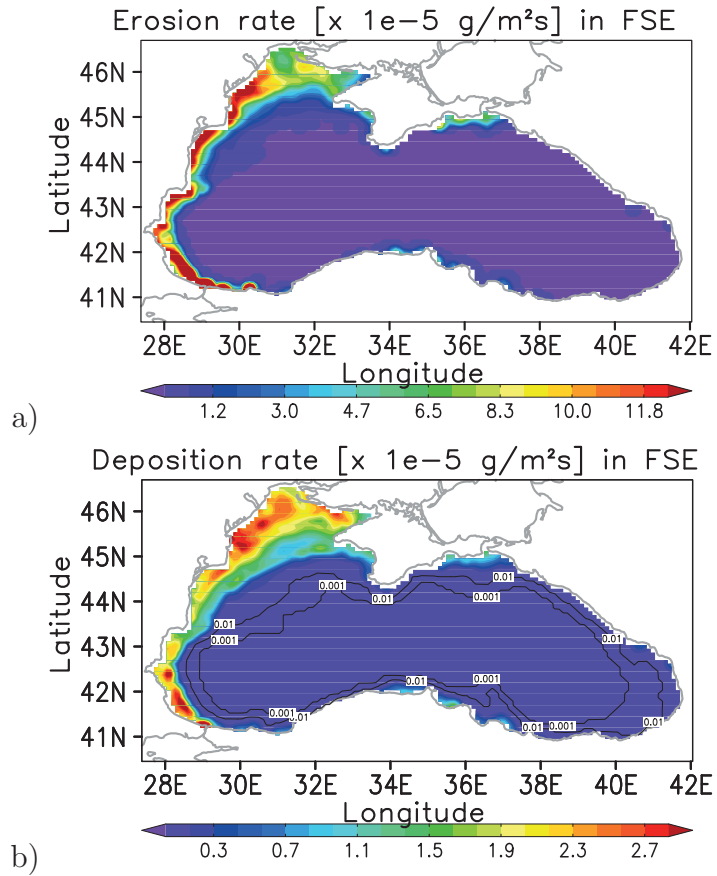


Figure 14: Annual mean erosion rate (a) and annual mean deposition rate (b) and as simulated in FSE.

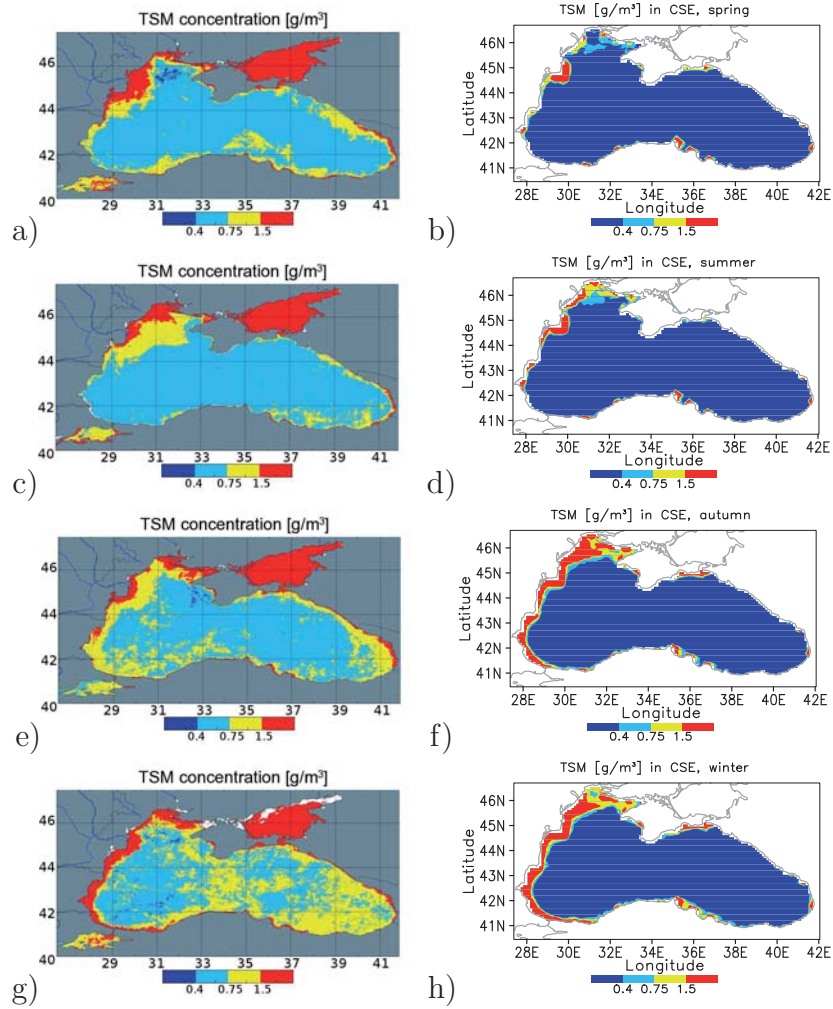


Figure 15: Seasonally averaged remote sensing TSM concentrations processed by ACRI ST(left column: a, c, e, and g) and model simulated seasonal mean TSM concentrations as in CSE (right column: b, d, f, and h). First row is spring, second one is summer, third one is fall, the last one is winter.

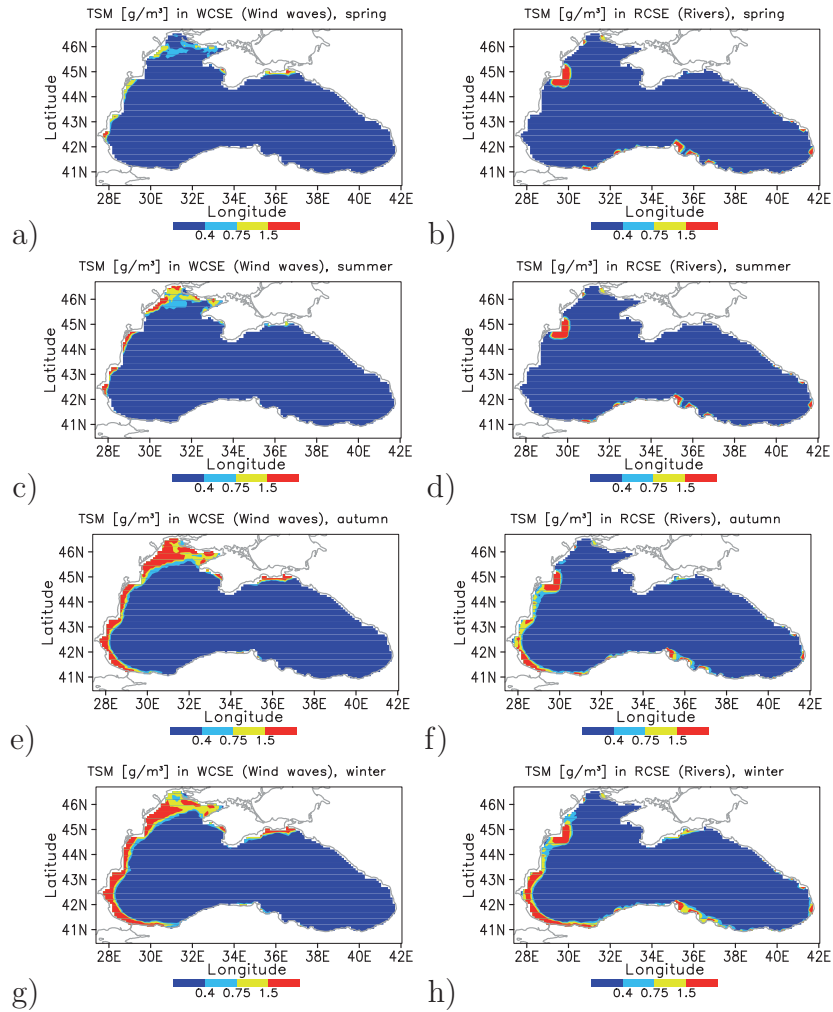


Figure 16: The impact of river and wind-wave forcing. Seasonally averaged TSM concentrations as simulated in WCSE and RCSE. The first column - WCSE(a, c, e, and g) corresponds to experiment with no river forcing (no sediment flux in the river mouth). In the second column - RCSE (b, d, f and h) the wind wave forcing has been neglected.